



Active region emergence and its effect on the solar corona



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Thanks

- George Fisher (UCB)
- Alex Pevtsov (NSO)

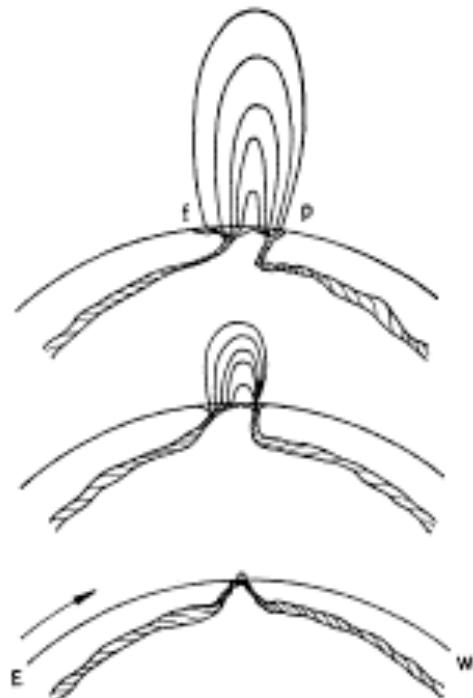
- Isaac Klapper
 - B. Ravindra*
 - Brian Welsch[§]
- } MSU

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* Presently IIA

§ Presently UCB

Active regions: where they come from

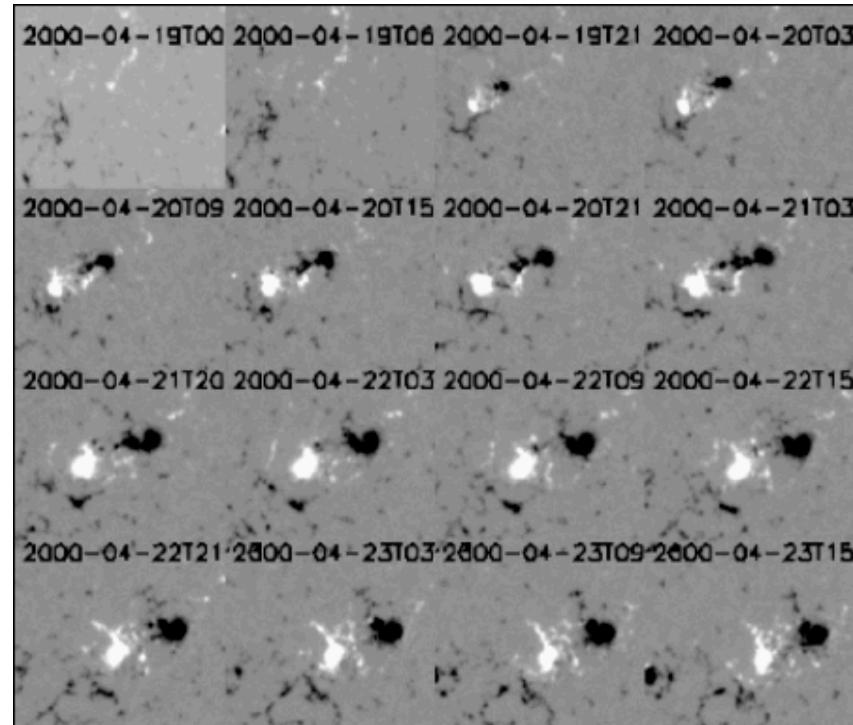


Babcock 1961

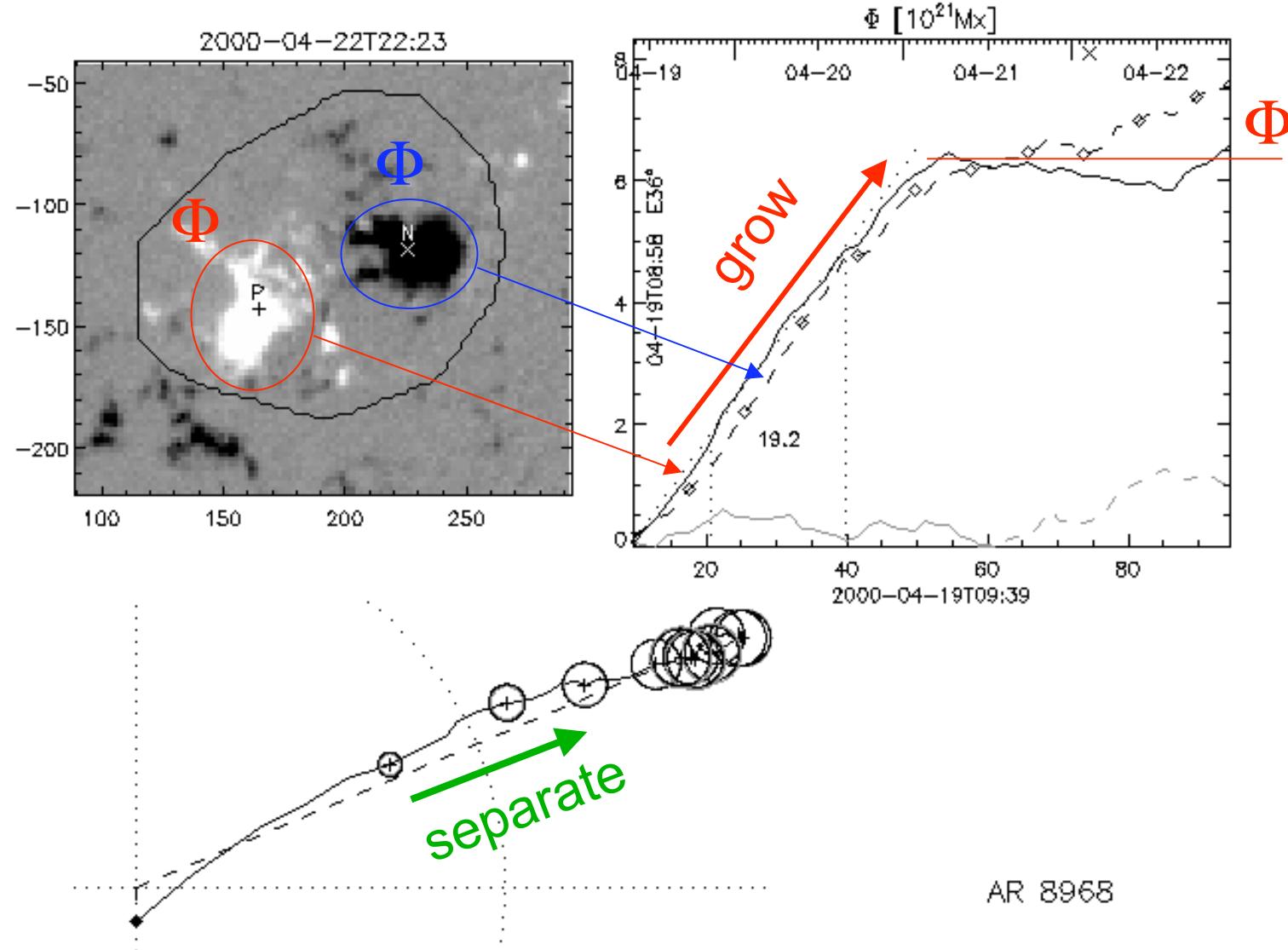
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Typical AR:
8968

[movie](#)



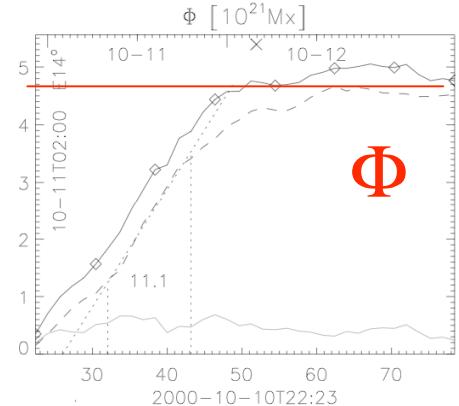
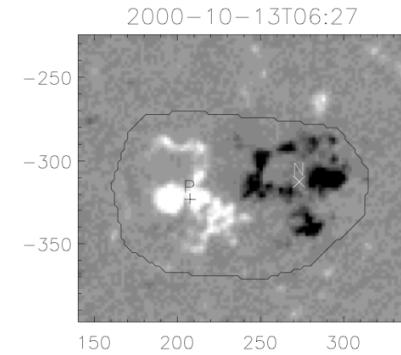
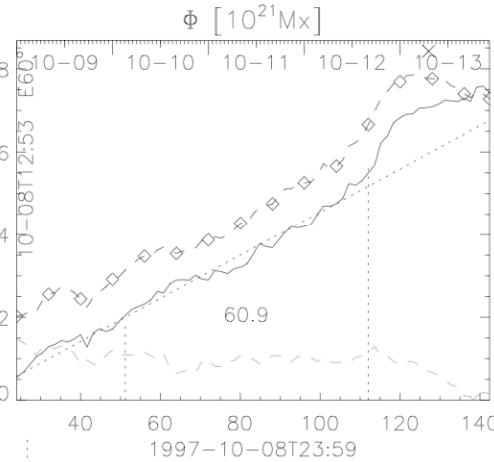
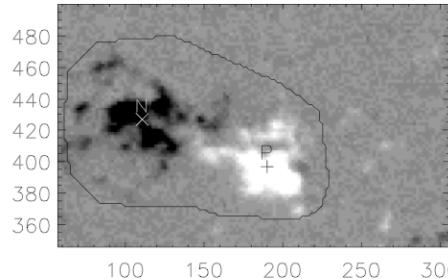
MDI



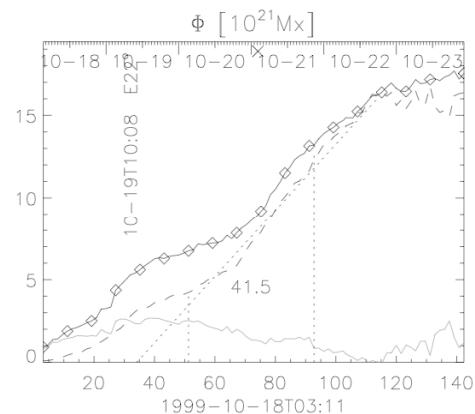
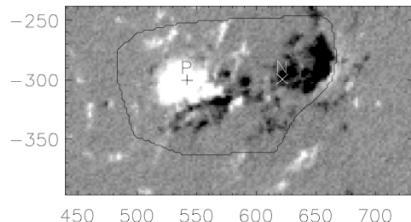
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various sizes...

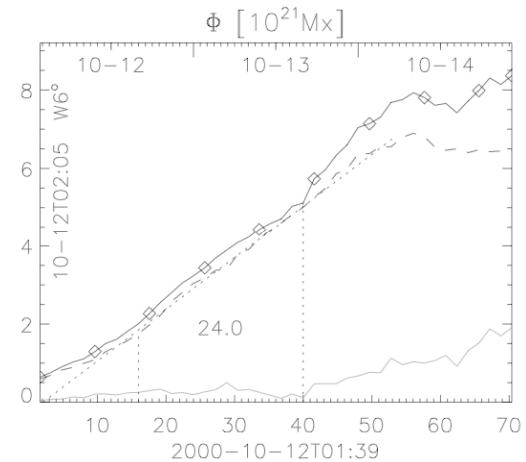
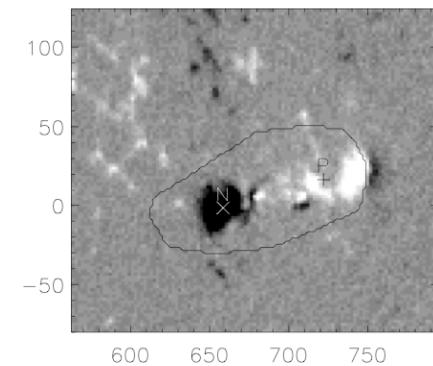
1997-10-13T22:24



1999-10-23T22:23

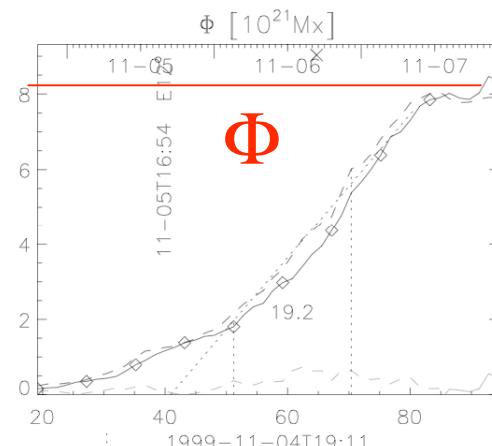
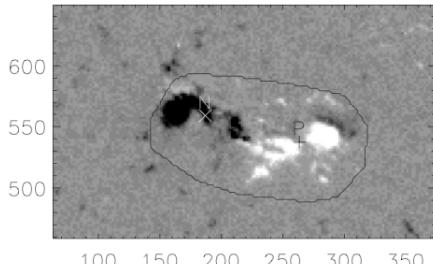


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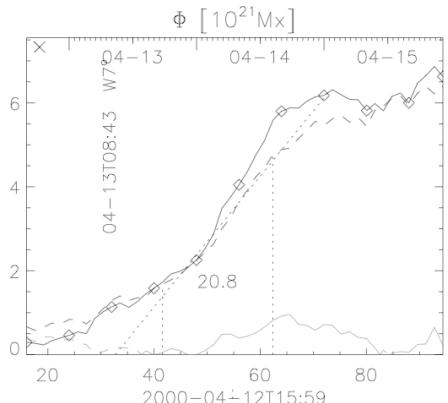
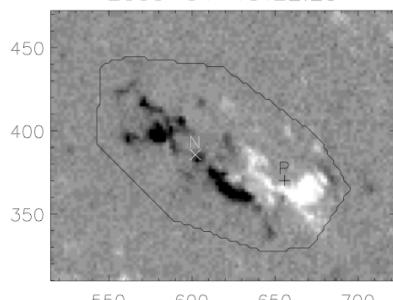


... same story

1999-11-07T22:24



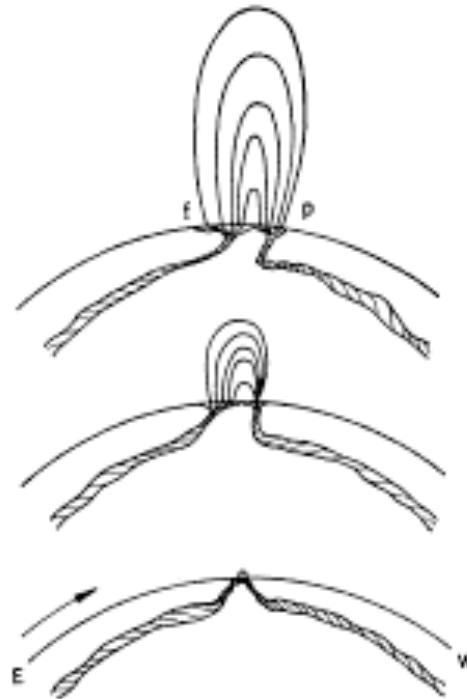
2000-04-15T22:23



How do these emerging flux tubes affect the corona?

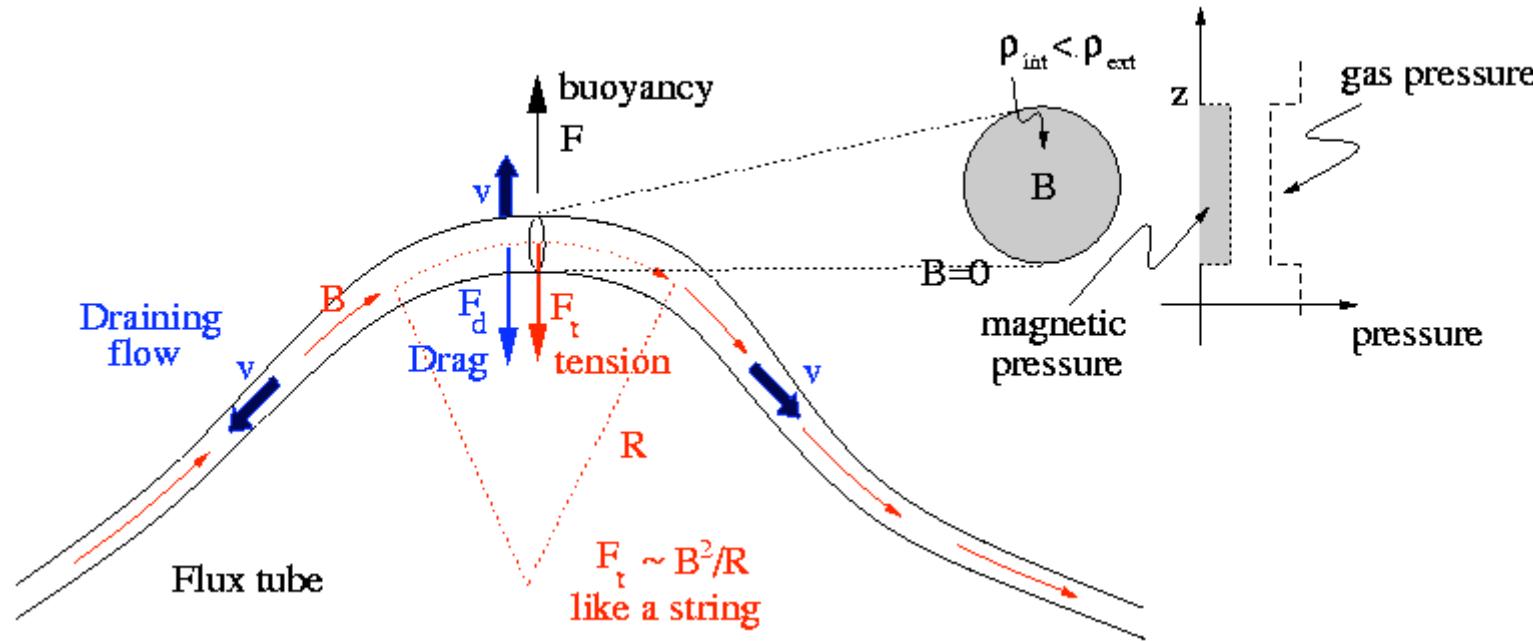
Outline

1. Dynamics of emergence
2. Twist (helicity) in emerging tubes
3. Transport of helicity into the corona by emerging tubes



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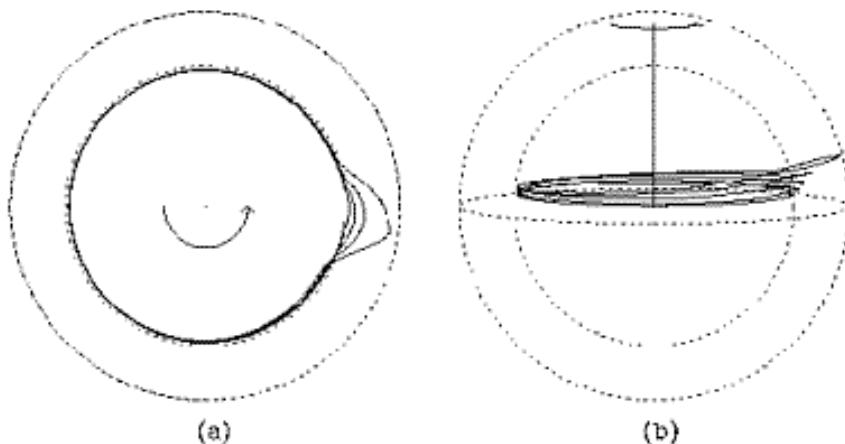
Dynamics of rising flux tubes



- Isolated tube, pressure-confined, “thin” $\leftrightarrow a \ll H_p$
- Axis of tube: space curve $\mathbf{x}(s,t)$
- Dynamical equations: Spruit 1981, Choudhuri & Gilman 1987

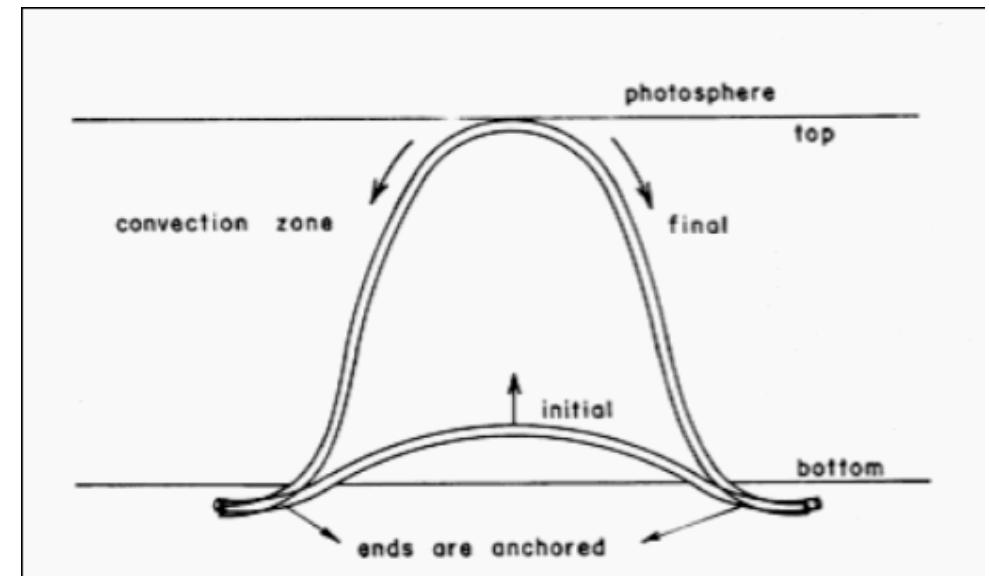
Model evolution of AR tubes

- Initialize tube at base of CZ
- Follow evolution of emerging tube - thin FT eqns.
- Predict configuration of observed AR



Fan *et al.* 1994

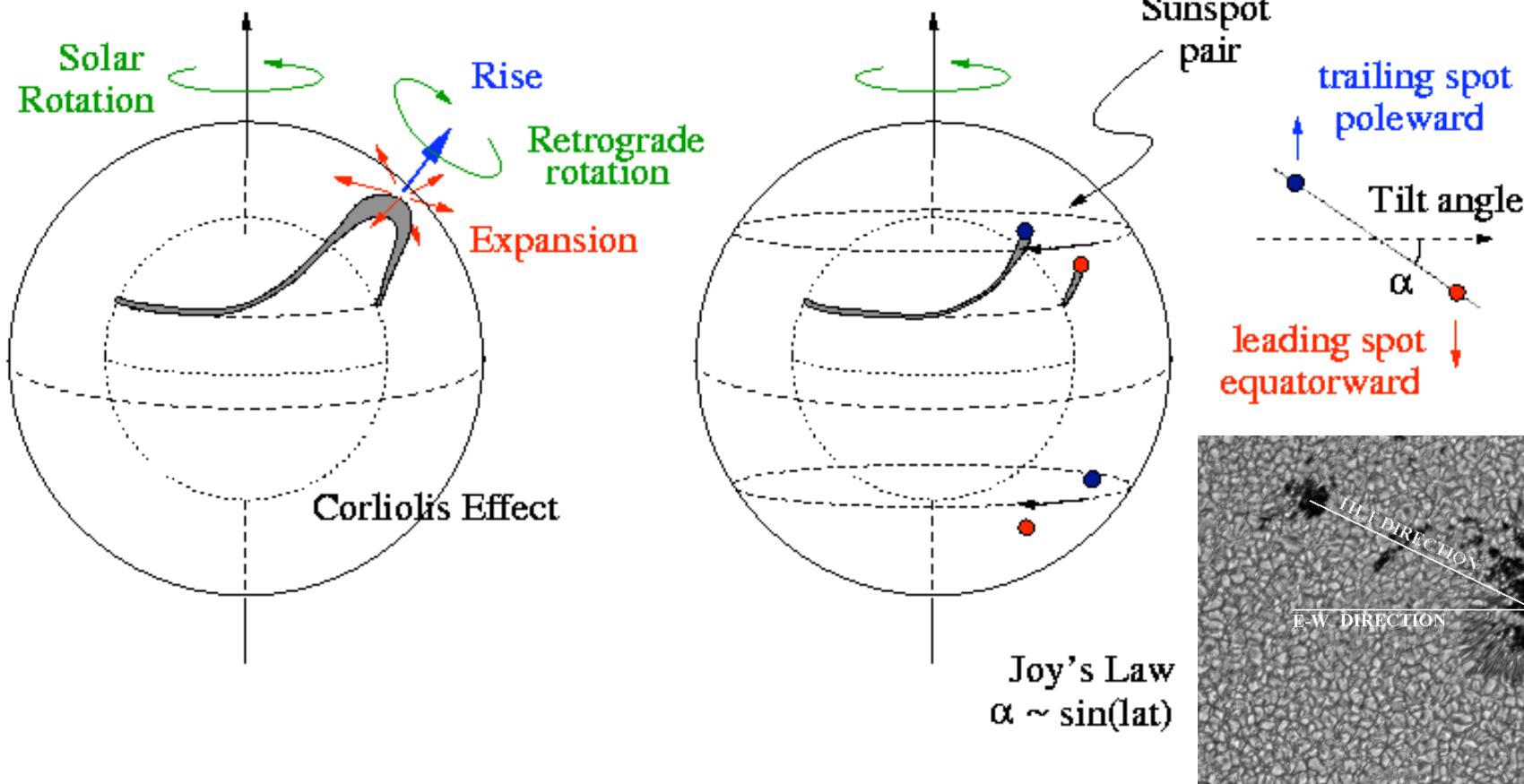
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D'Silva & Choudhuri 1993

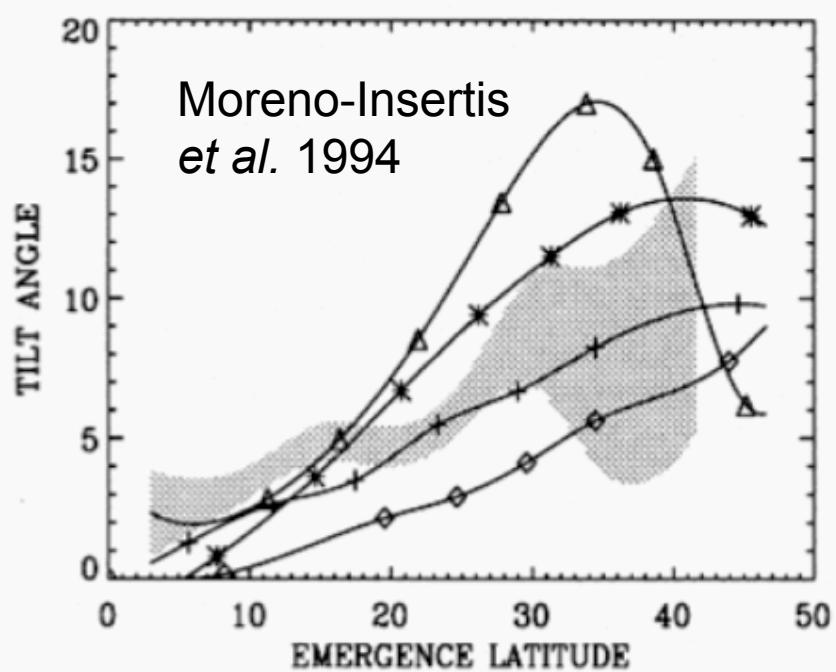
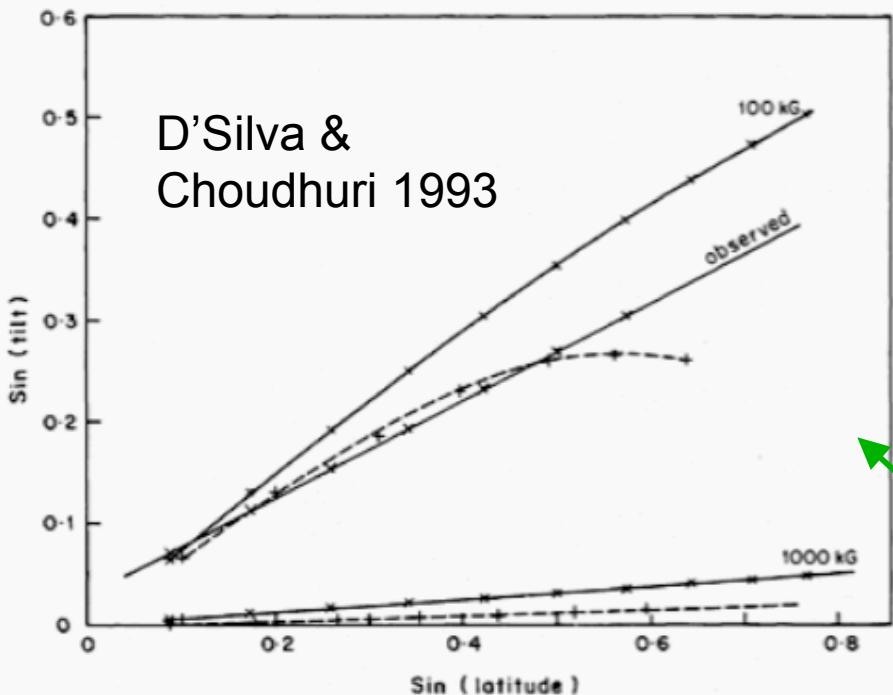
A Rising Flux Tube

Deflection of rising tube by Coriolis effect → tilted pair of spots



Thin flux tube successes:

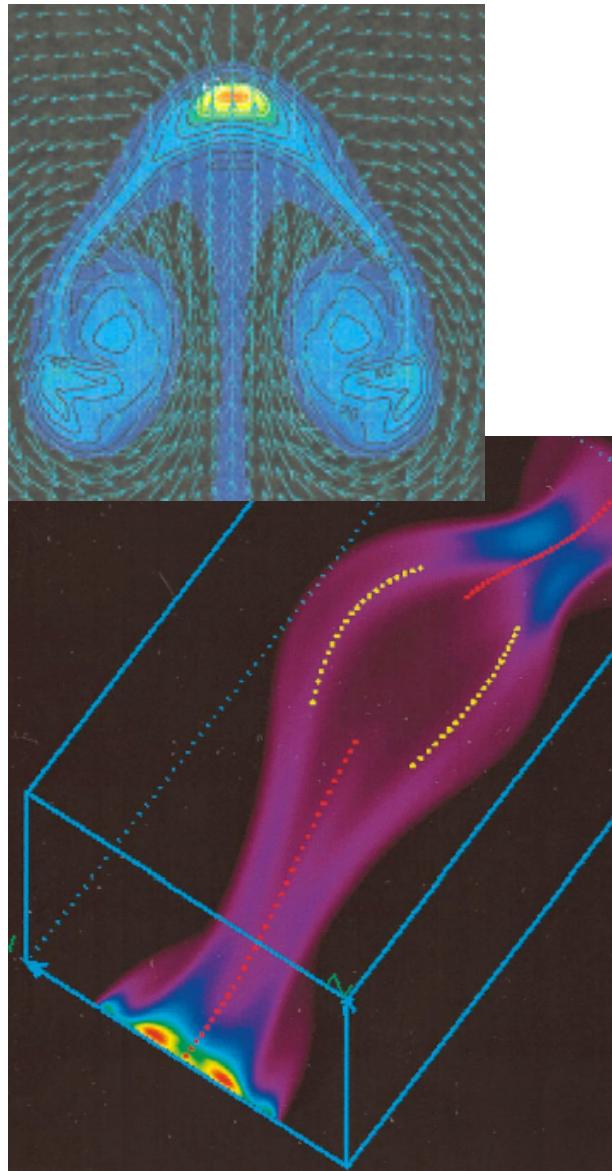
- Hale's polarity law
- Joy's law for tilt angles
(D'Silva & Choudhuri 1993)
- p-f asymmetry
(Fan *et al.* 1993)
- post-emergence velocities
(Moreno-Insertis *et al* 1994)
- Statistical dispersion
(Longcope & Fisher 1996)



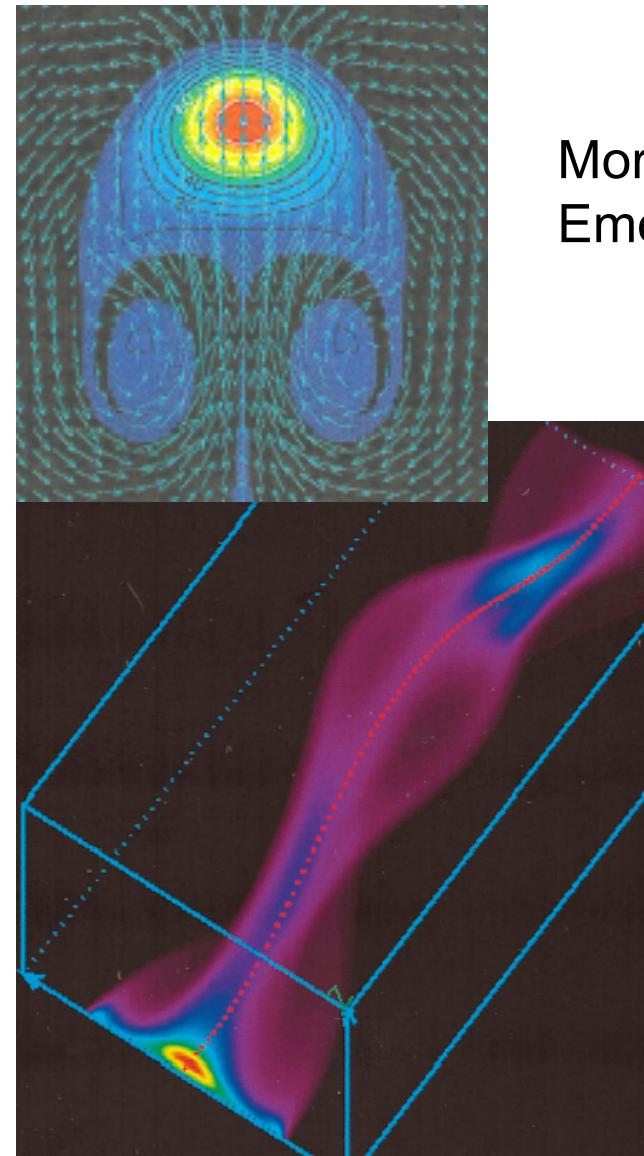
Flux Tube Twist

Flux tubes must be twisted in order to rise (Parker 1979)

untwisted



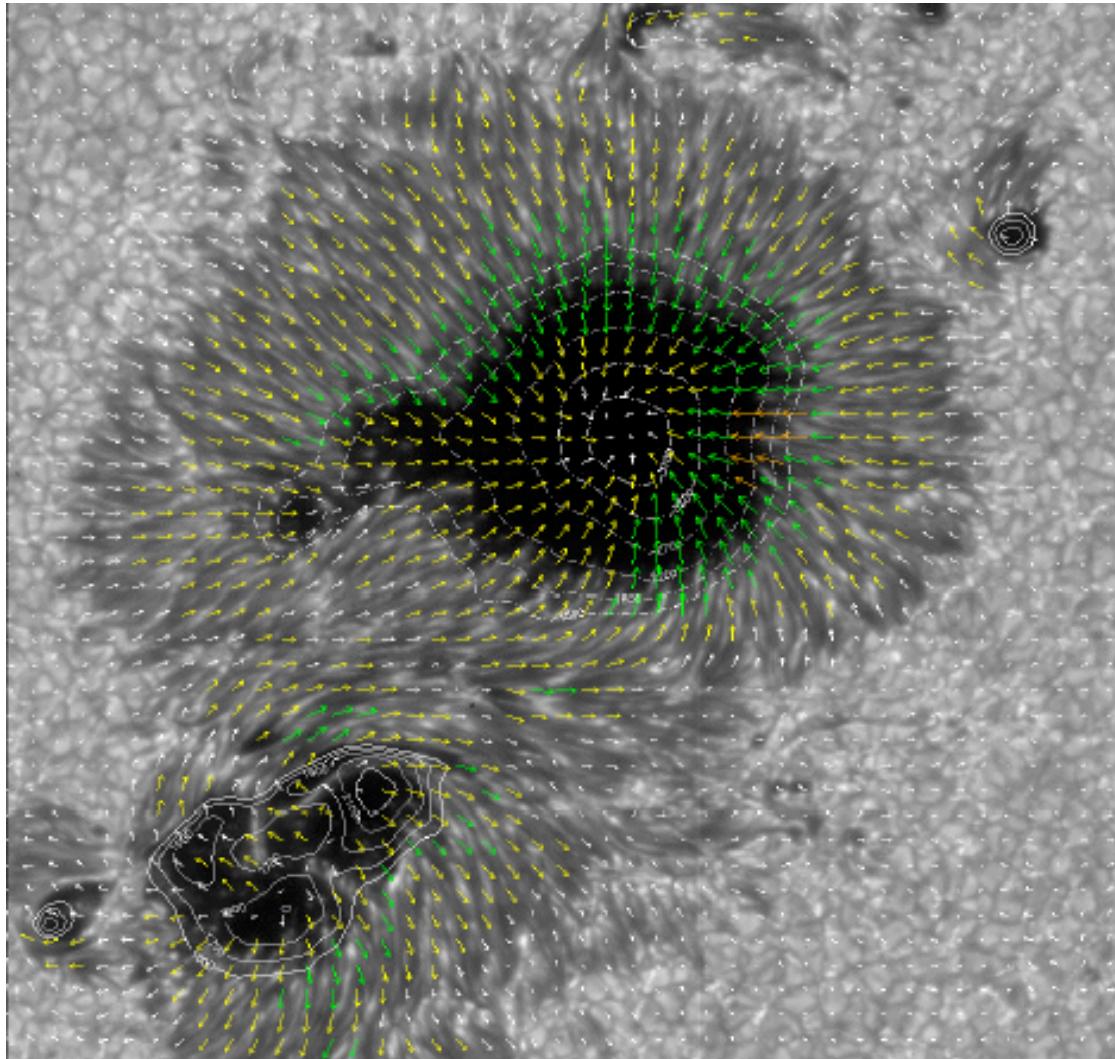
twisted



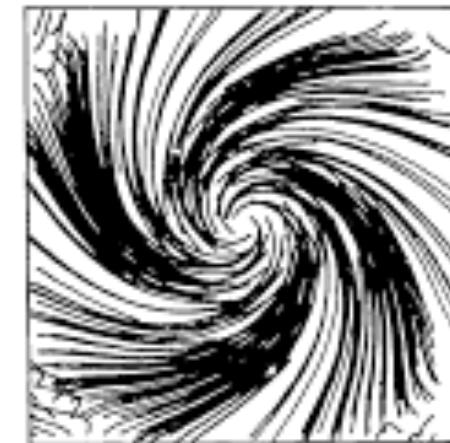
Moreno-Insertis &
Emonet 1996

Abbott et al. 2000

... and AR fields are twisted



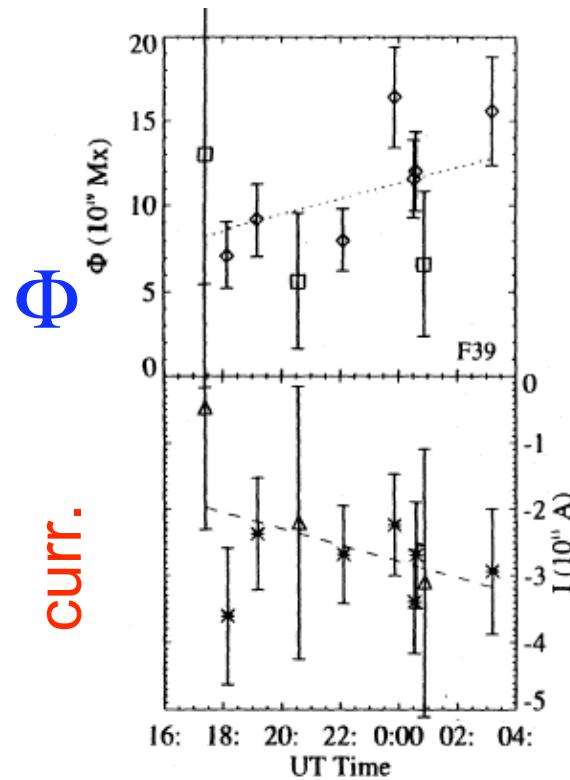
(courtesy T. Magara & Hinode)
Danyalov, India



b)

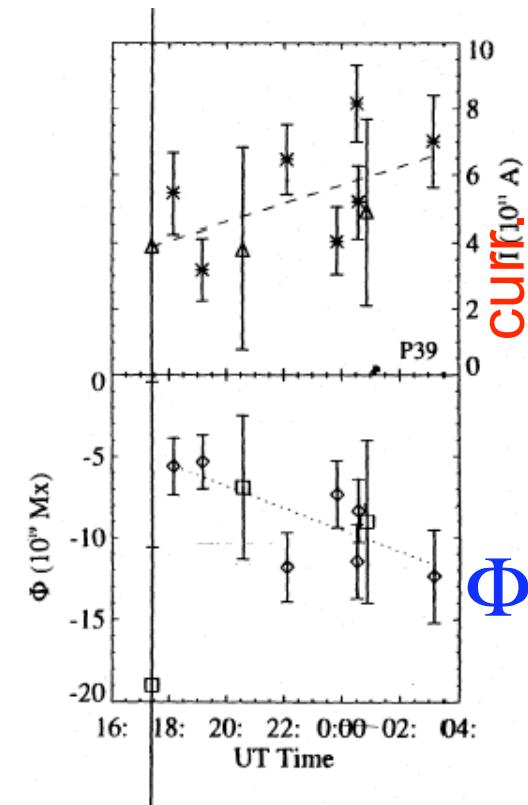
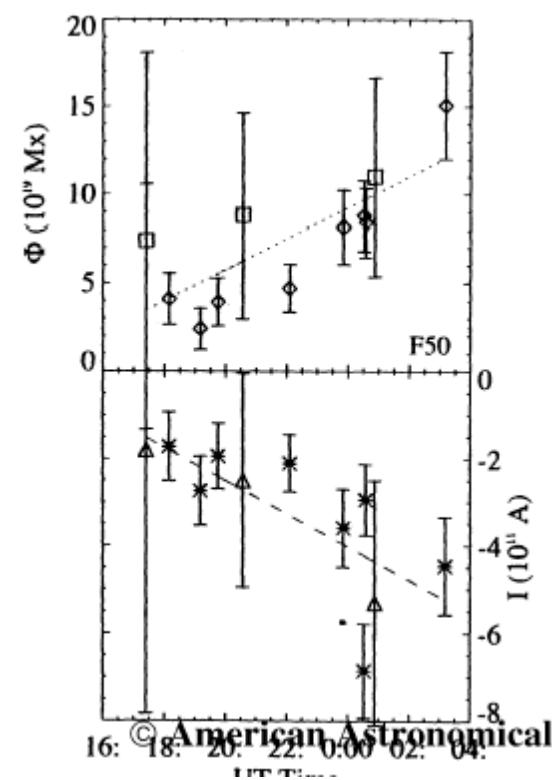
(from Nakagawa
et al. 1971)

Evidence that flux tubes emerge already twisted: Flux (Φ) and current increase together (Leka *et al.* 1996)

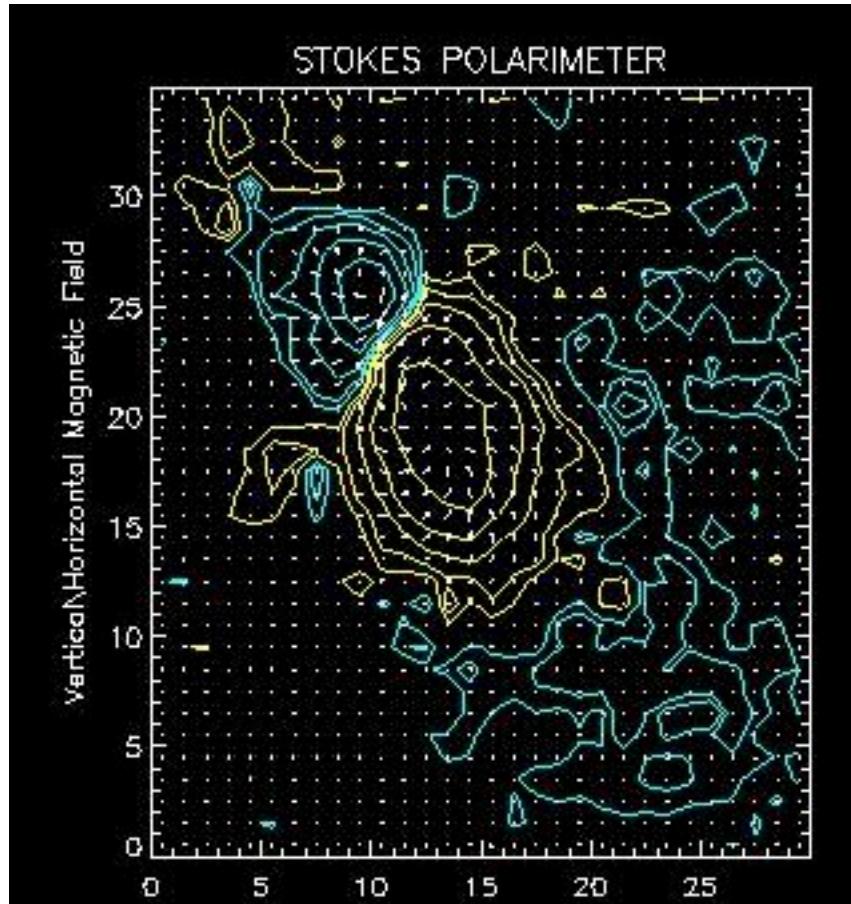


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How twisted are the tubes?



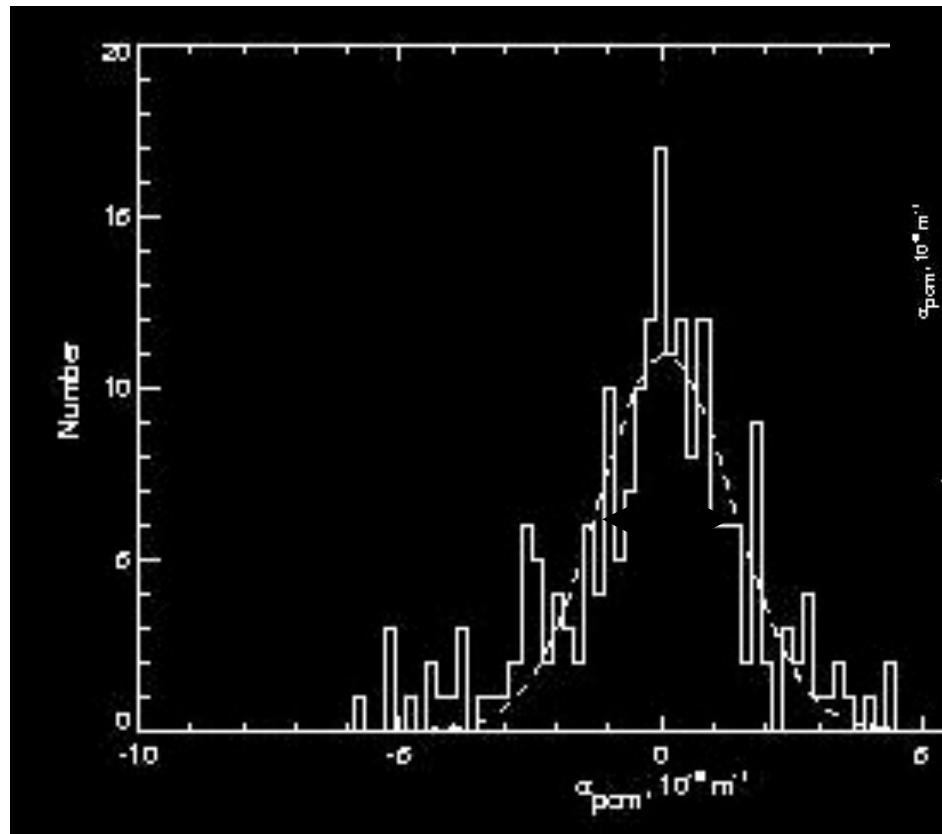
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α_{best} introduced by
Pevtsov, Canfield
& Metcalf (1995)

- calc $\mathbf{B}_{\perp}^{(\alpha)}$ by extrapolating B_z w/ fixed value of α
- vary α until $\mathbf{B}_{\perp}^{(\alpha)}$ best matches observed \mathbf{B}_{\perp}
i.e. minimize

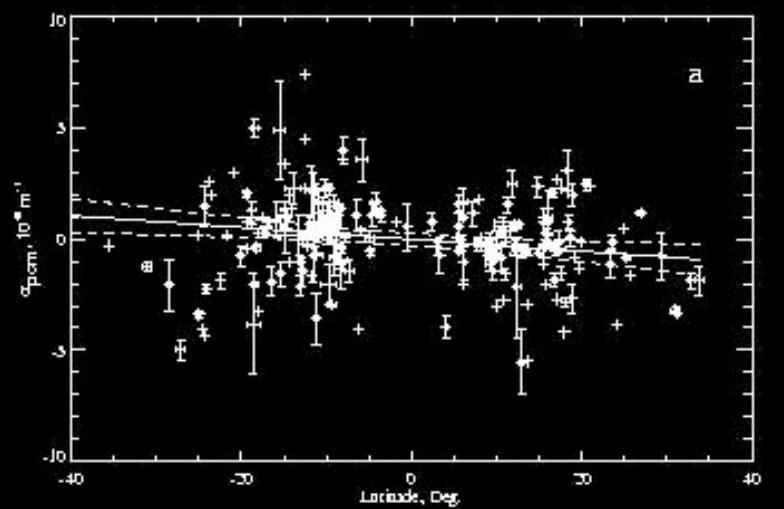
$$\sum_{i,j} |\mathbf{B}_{\perp}(i,j) - \mathbf{B}_{\perp}^{(\alpha)}(i,j)|^2$$

How twisted are the tubes?



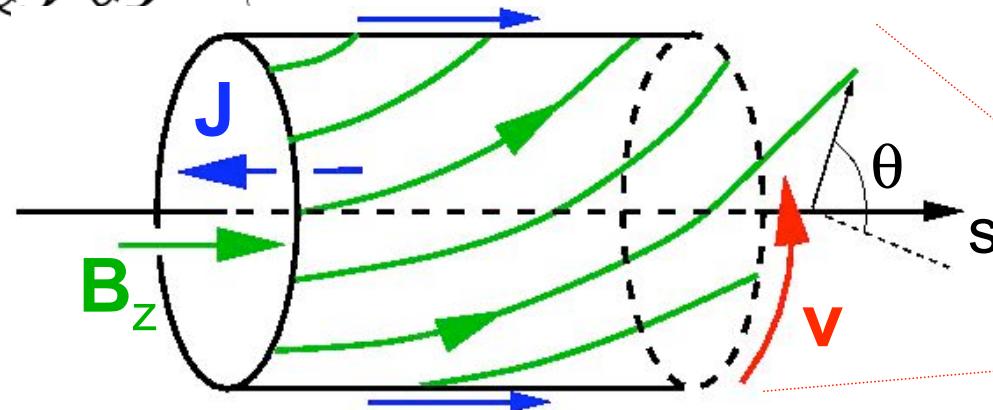
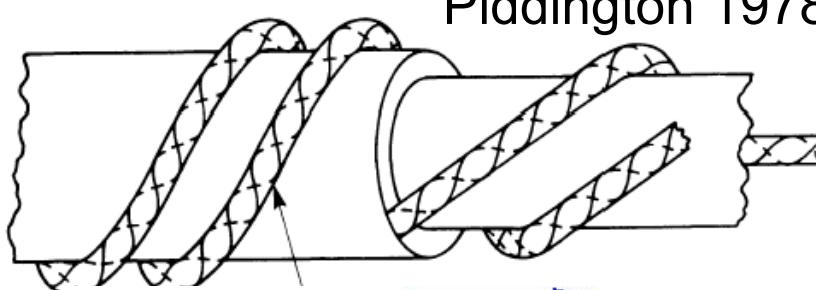
Linear trend removed
(from Longcope, Fisher &
Pevtsov 1998)

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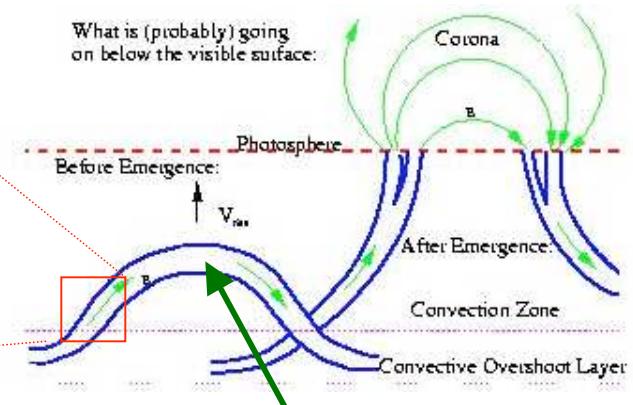


$\langle \alpha_{\text{best}} \rangle \sim 3 \times 10^{-9} \text{ m}^{-1}$
varies /w latitude

$\Delta \alpha_{\text{best}} \sim 10^{-8} \text{ m}^{-1}$
independent of latitude



Twist in flux tubes



Field lines twist about axis at a rate

$$q(s,t) \text{ “=} \frac{d\theta}{ds}$$

$$\alpha = 2q$$

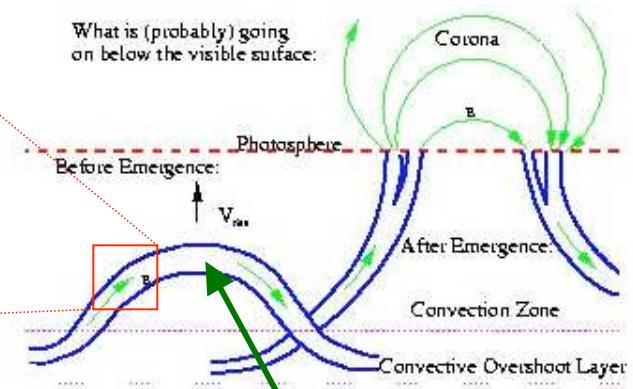
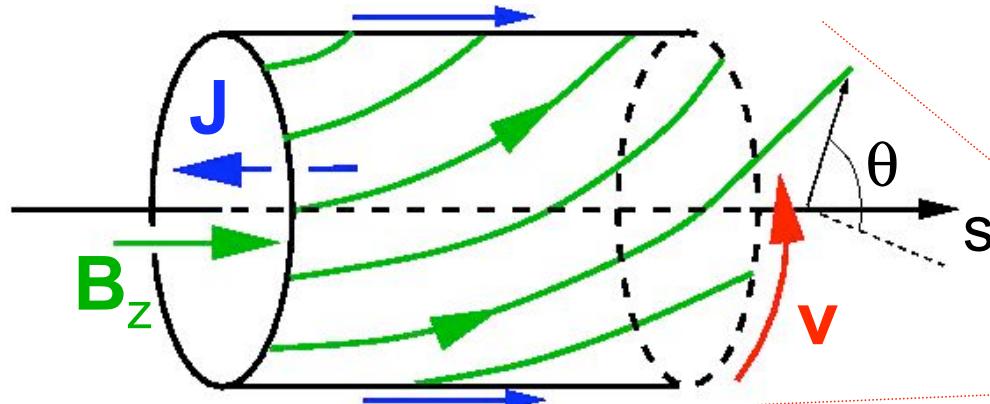
$$B_\theta = q r B_z$$

Plasma spins about axis at rate

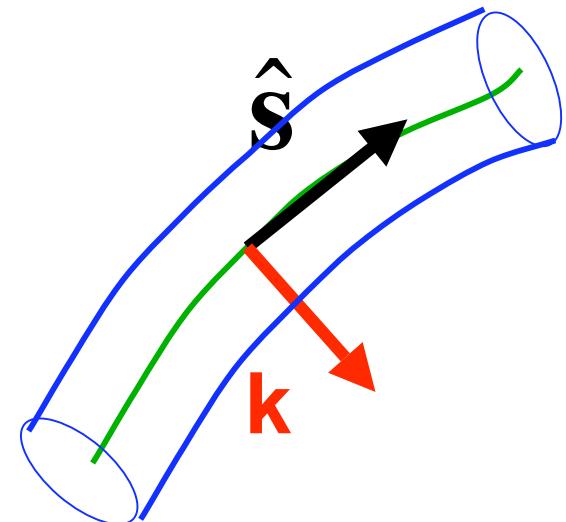
$$\omega(s,t) \text{ “=} \frac{d\theta}{dt}$$

$$v_\theta = \omega r$$

Twist in flux tubes



Axis of tube:

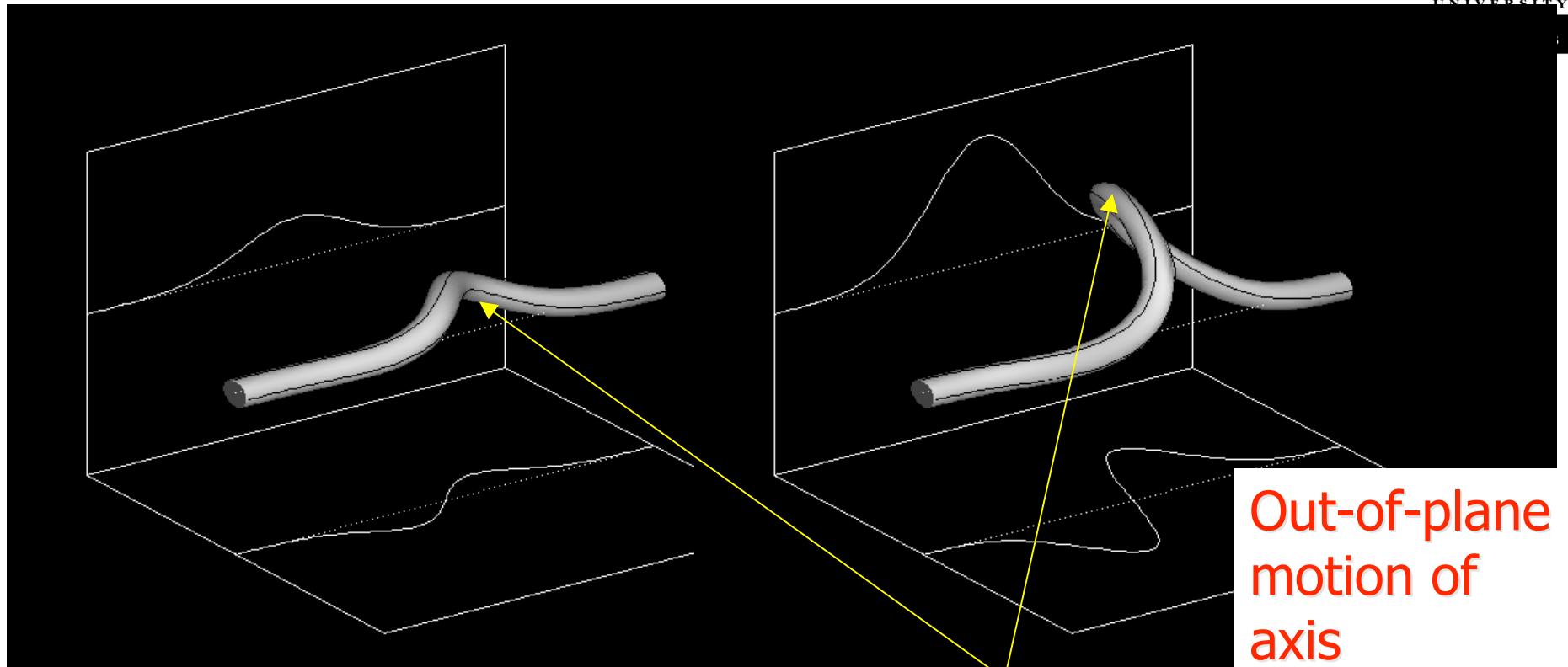


Evolution of twist & spin

$$\frac{d\omega}{dt} = v_A^2 \frac{\partial q}{\partial s} - \frac{2}{a} \frac{da}{dt}$$

$$\frac{dq}{dt} = \frac{\partial \omega}{\partial s} - \left(\hat{s} \cdot \frac{\partial \mathbf{v}}{\partial s} \right) q + (\hat{s} \times \mathbf{k}) \cdot \frac{\partial \mathbf{v}}{\partial s}$$

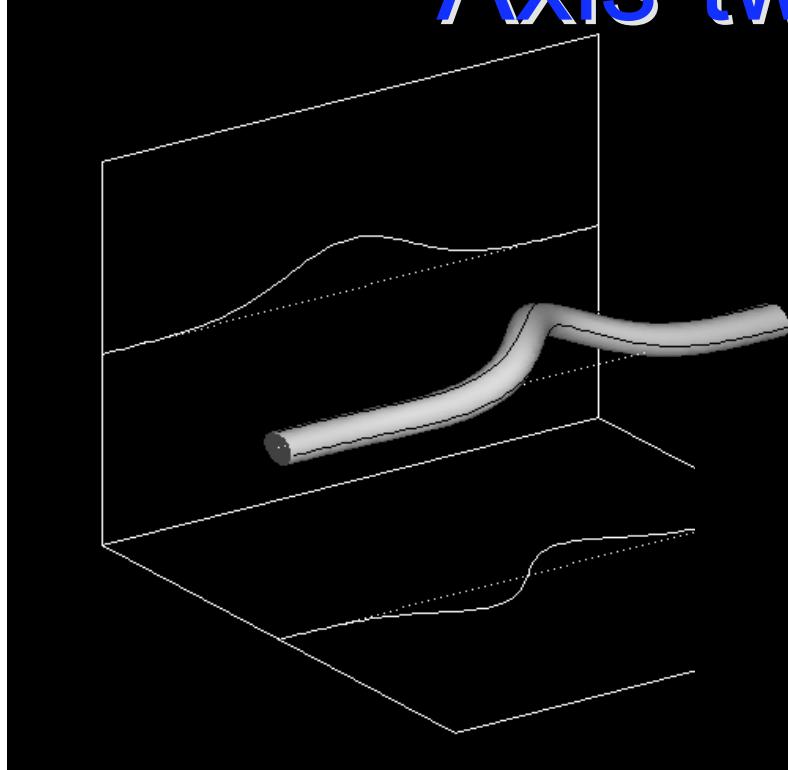
Dynamics of twist



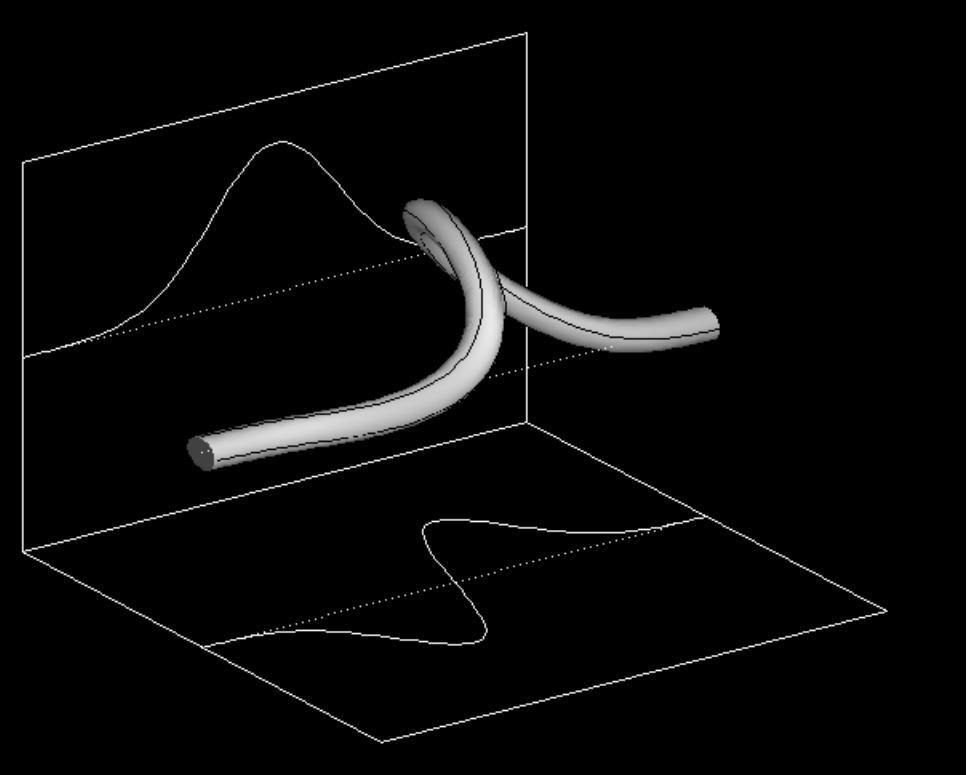
$$\frac{dq}{dt} = \frac{\partial \omega}{\partial s} - \left(\hat{s} \cdot \frac{\partial \mathbf{v}}{\partial s} \right) q + \left(\hat{s} \times \mathbf{k} \right) \cdot \frac{\partial \mathbf{v}}{\partial s} \xleftarrow{\Sigma(s)}$$

indep. of q or ω

Axis-twist coupling



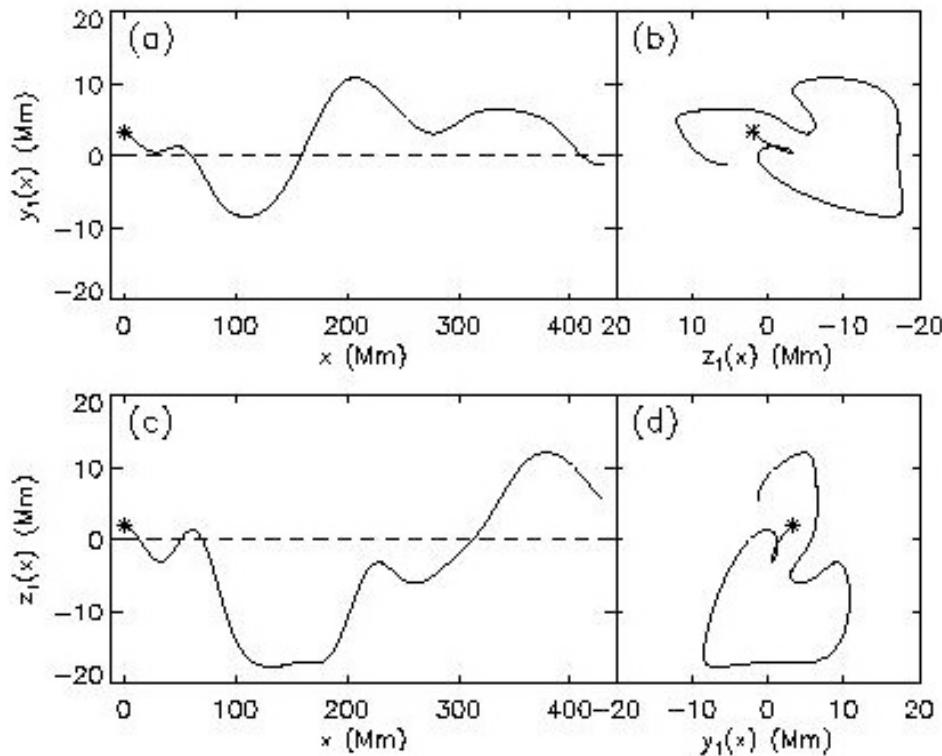
$$\frac{dq}{dt} = \Sigma = (\hat{\mathbf{s}} \times \mathbf{k}) \cdot \frac{\partial \mathbf{v}}{\partial s} > 0$$



Increasing LH
writhe ($dW_r/dt < 0$)
 → Increasing RH
twist ($dTw/dt > 0$)

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Writhe from Turbulence: The Σ -effect (Longcope, Fisher & Pevtsov 1998)



Spectrum of kinetic helicity

$$\langle \mathbf{v}^e \cdot \nabla \times \mathbf{v}^e \rangle = -\frac{1}{3} \tau_c \int F(k) dk$$

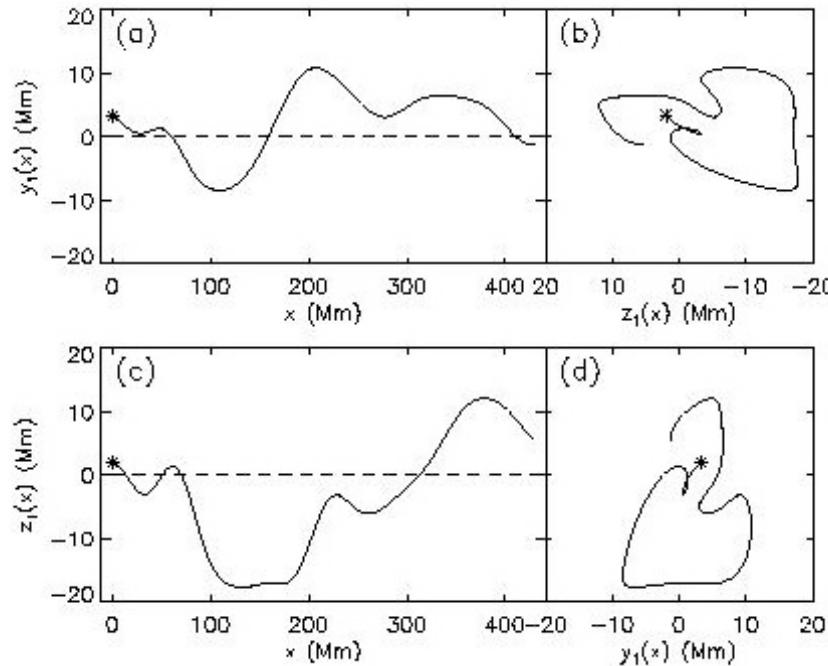
$$\Sigma \approx (\hat{\mathbf{s}} \times \mathbf{k}) \cdot \frac{\partial \mathbf{v}^e}{\partial s} \quad \text{Twist source}$$

Averaging over turbulence:

$$\langle \Sigma \rangle \approx -\frac{1}{5} \tau_c \int k^2 F(k) dk$$

Variance of twist source: $\langle \Sigma^2 \rangle \approx |\mathbf{k}|^2 \int k^2 E(k) dk \gg \langle \Sigma \rangle^2$

Σ -effect vs. α -effect



$$\langle \Sigma \rangle \approx -\frac{1}{5} \tau_c \int k^2 F(k) dk$$

Spectrum of kinetic helicity

Compare to α -effect:

$$\alpha \approx -\langle \mathbf{v}^e \cdot \nabla \times \mathbf{v}^e \rangle = \frac{1}{3} \tau_c \int F(k) dk$$

Cause of the observed twist



Observed properties

- Twist in CZ flux tube
- LH twist in North
- $\langle \alpha \rangle \sim 3 \times 10^{-9} \text{ m}^{-1}$
- 25% violation of trend
- $\Delta\alpha \sim 10^{-8} \text{ m}^{-1}$
- $\Delta\alpha$ indep. of latitude

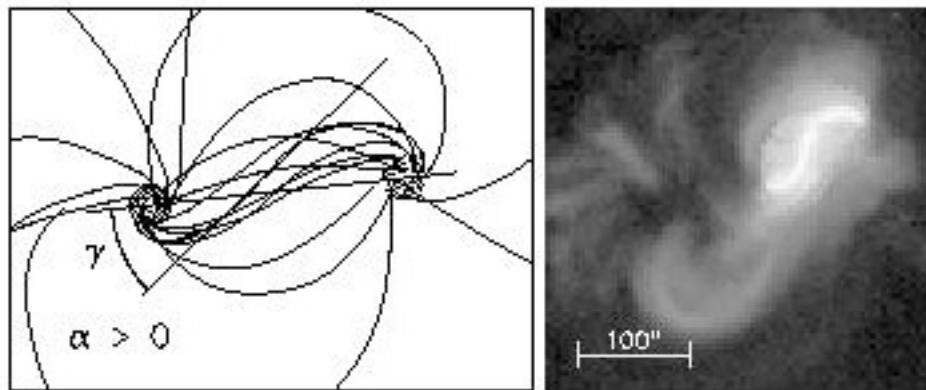


Kinetic helicity:

- RH writhe in North
- $\langle \alpha \rangle \sim 3 \times 10^{-9} \text{ m}^{-1}$
- Fluctuates (turbulence)
- Level indep. of latitude
- $\Delta\alpha \sim 10^{-8} \text{ m}^{-1}$

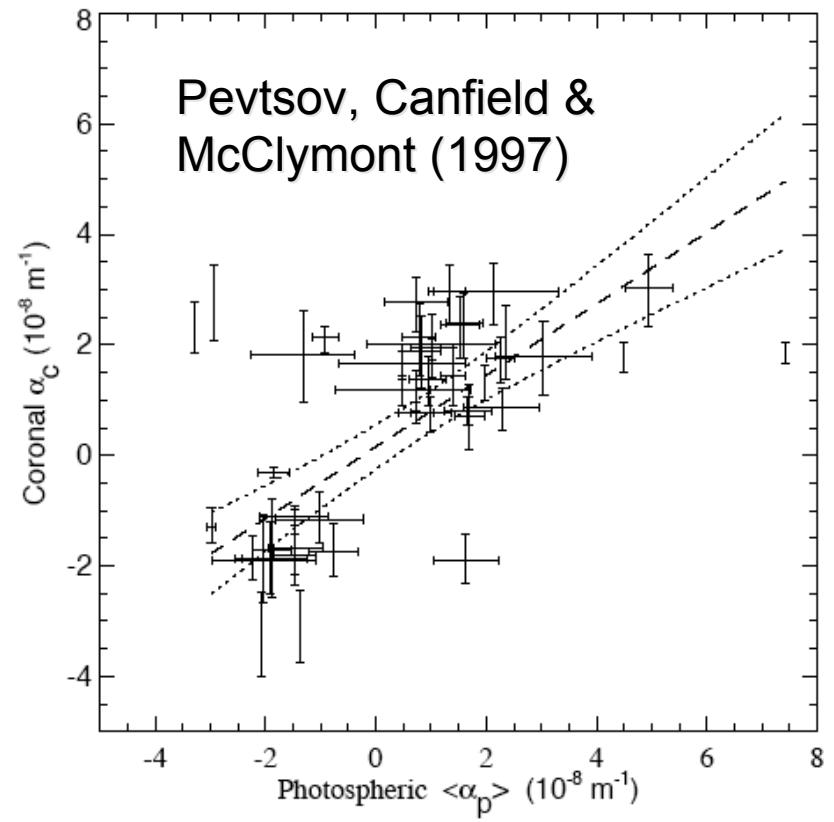
Twist: Photosphere vs. Corona

Force-free-field
w/ constant- α

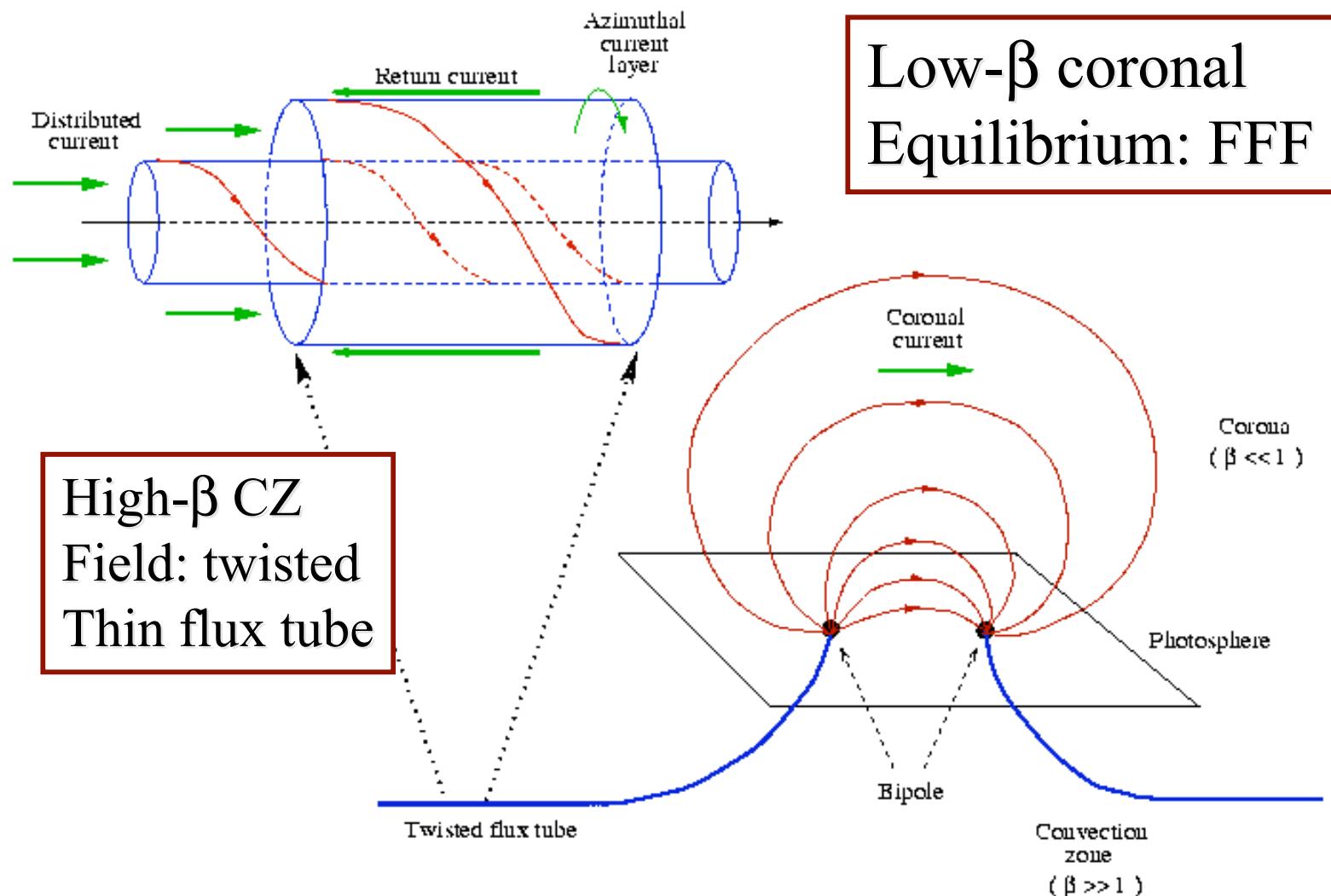


$\alpha_{\text{best}} \sim \alpha$ in coronal field

- α_{best} and α for 140 ARs
- Found α_{best} correlated with α



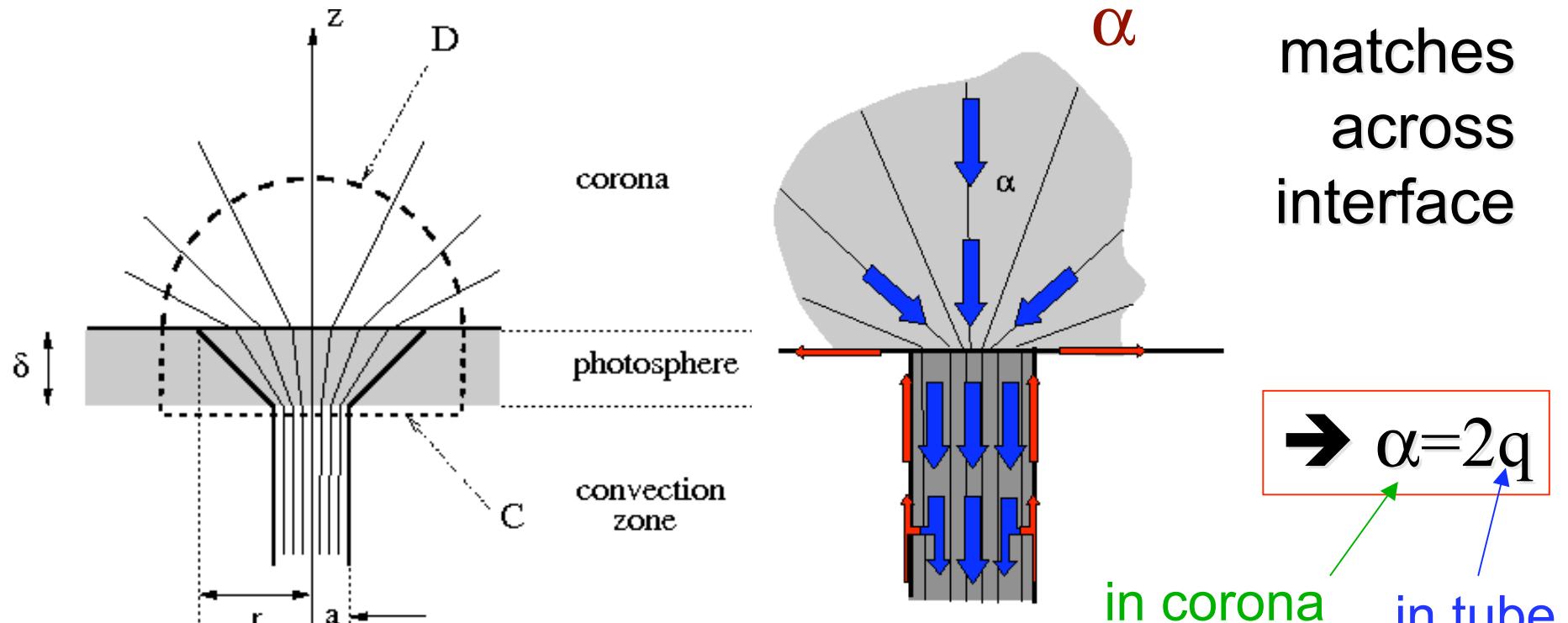
Coupling flux tube to corona



Coupling flux tube to corona



Balance of net torque



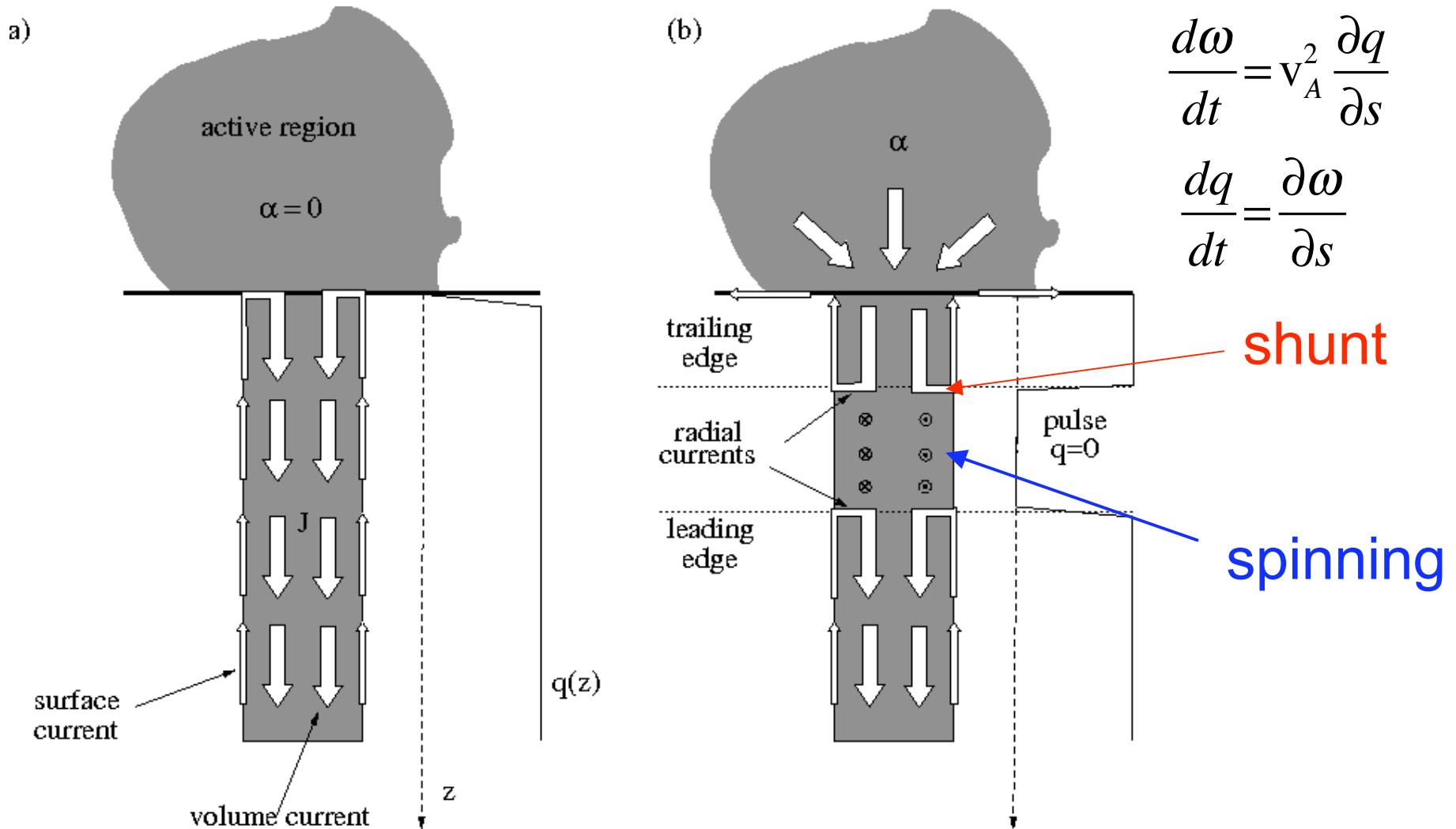
(Longcope & Weslch 1998)

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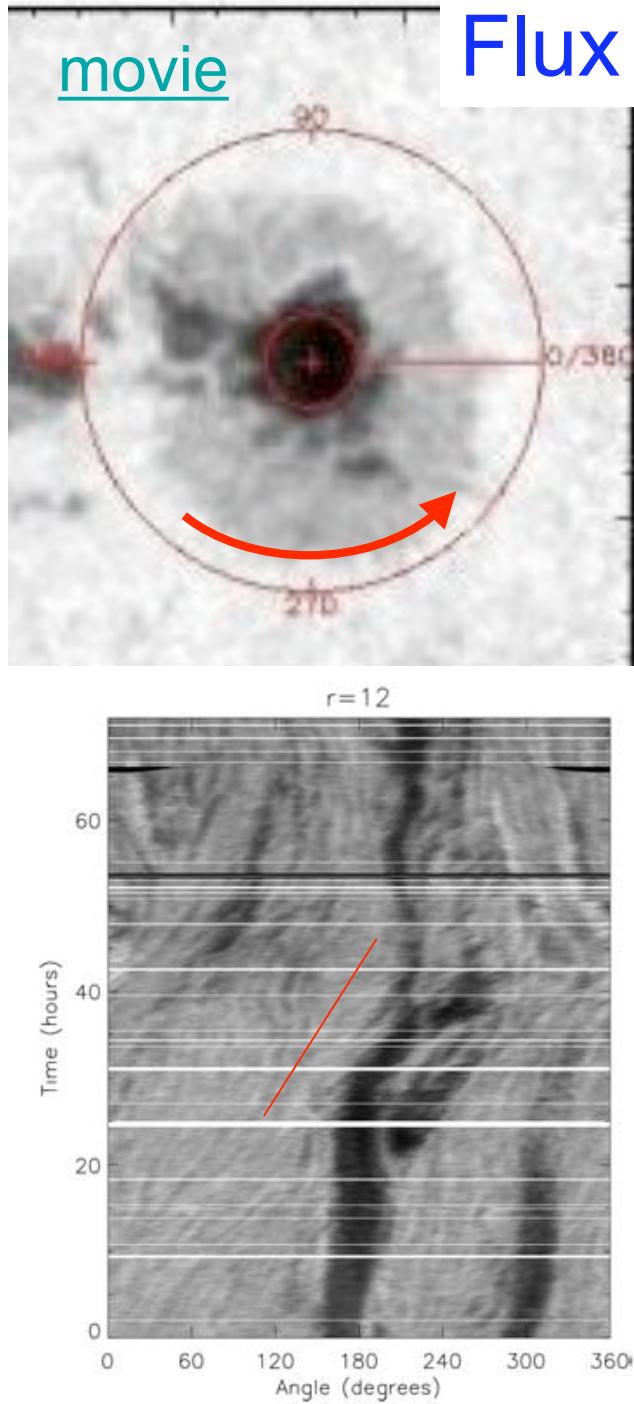
Coupling flux tube to corona



Imbalanced torque (shunted current) \rightarrow spin



movie

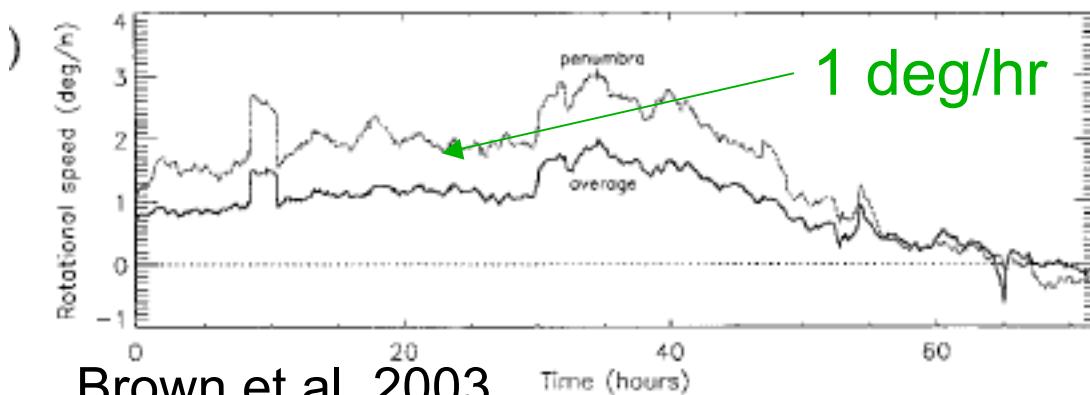


Flux tube twist → sunspot rotation

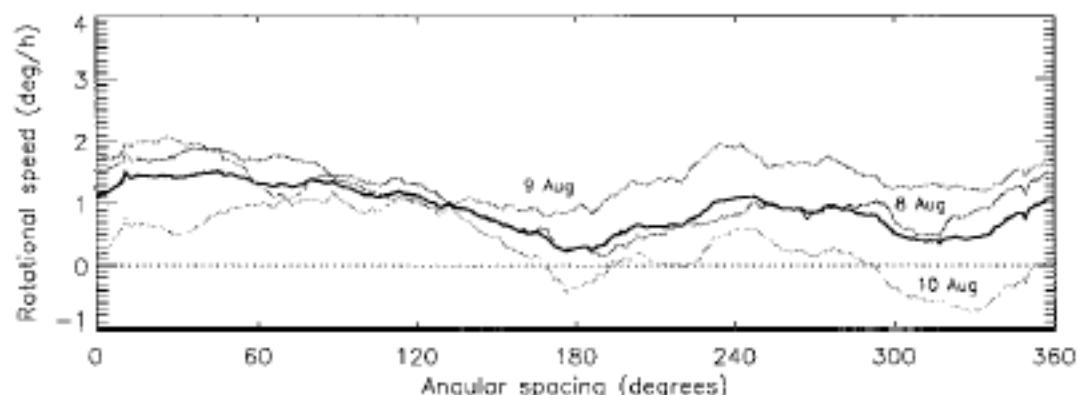
Summarising these results, we have a mean rotational movement of 0.35 km. per sec., and a radial motion of about 2 km. per sec. Combining them, we get a spiral or catharine-wheel movement opposite in the two hemispheres. This is represented graphically in the appended diagram (fig. 2), which shows the directions of the stream-lines in a northern and in a southern spot.



Evershed 1910

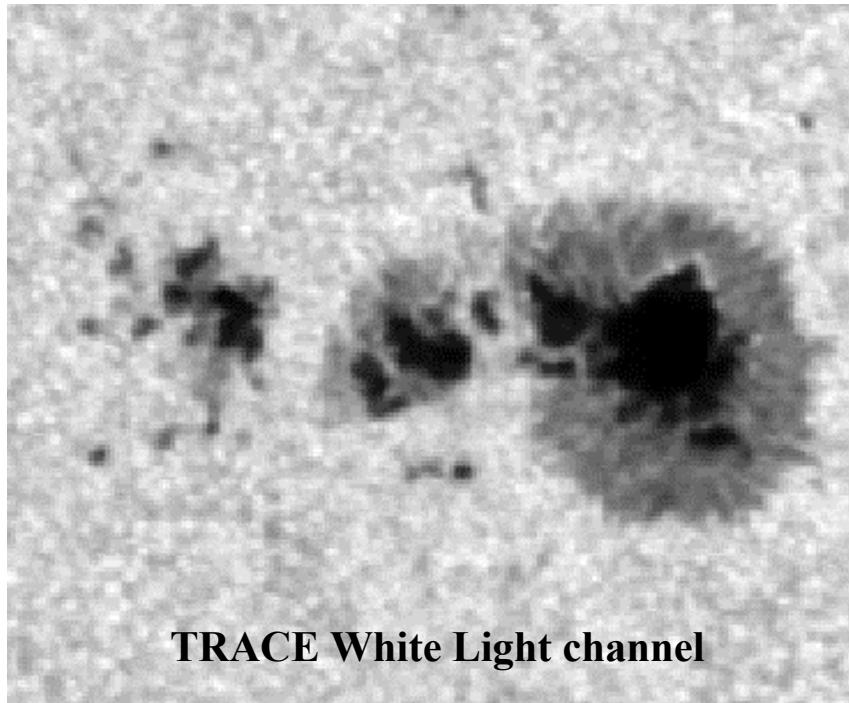


Brown et al. 2003

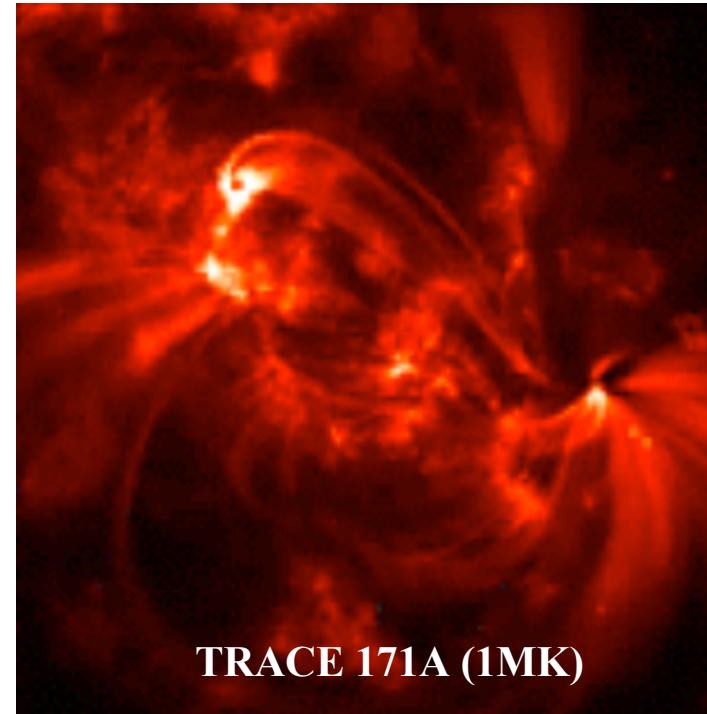




Twist Creates Spin



TRACE White Light channel



TRACE 171A (1MK)

movie

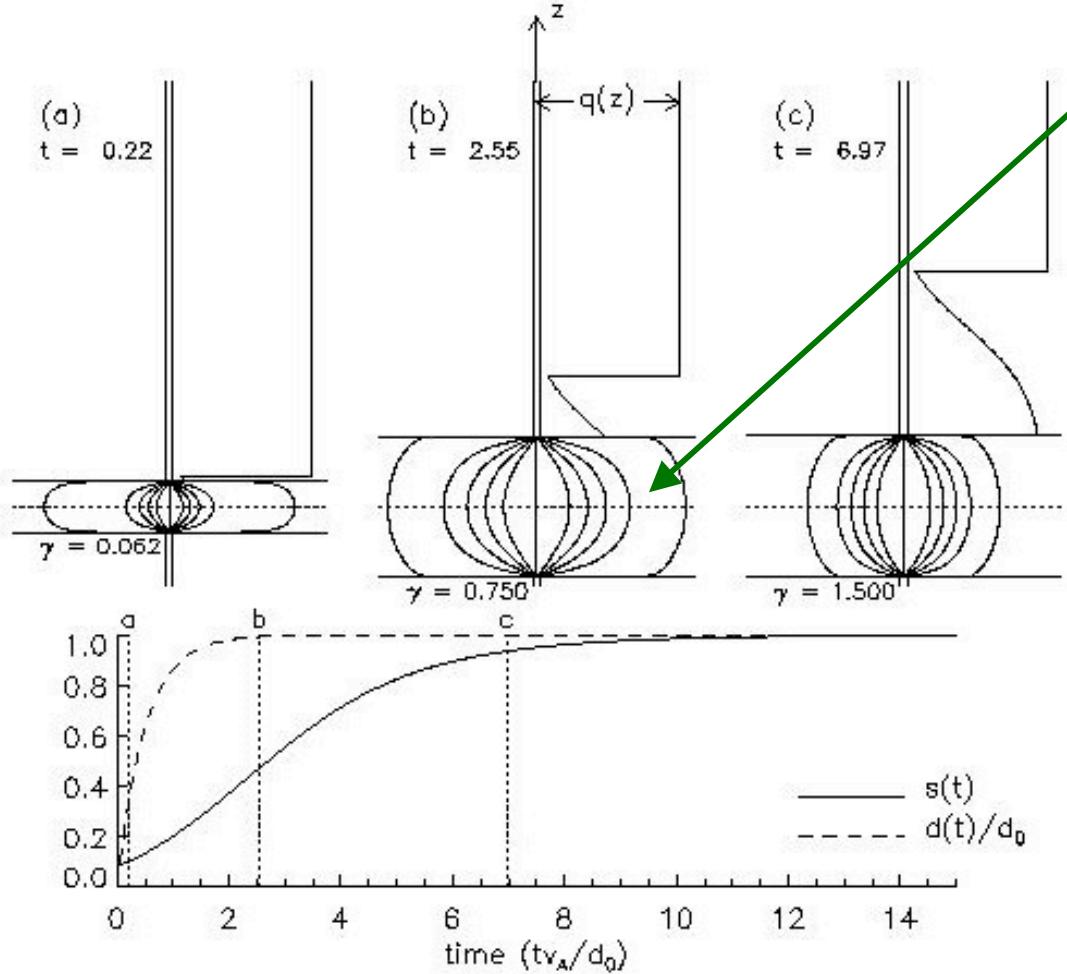
(Courtesy D. Alexander)

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Spin from Emergence

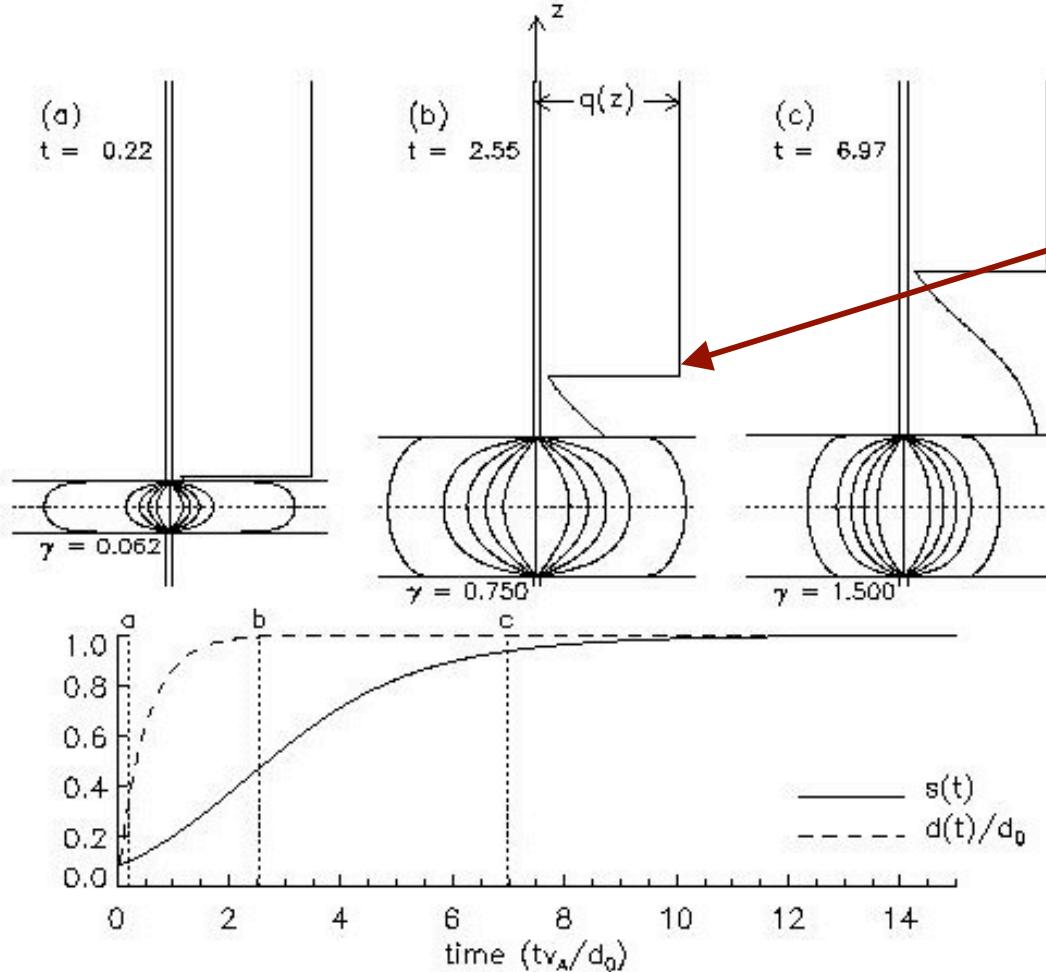
simple model: Longcope & Welsch 1998

- Twist propagates into corona



Spin from Emergence

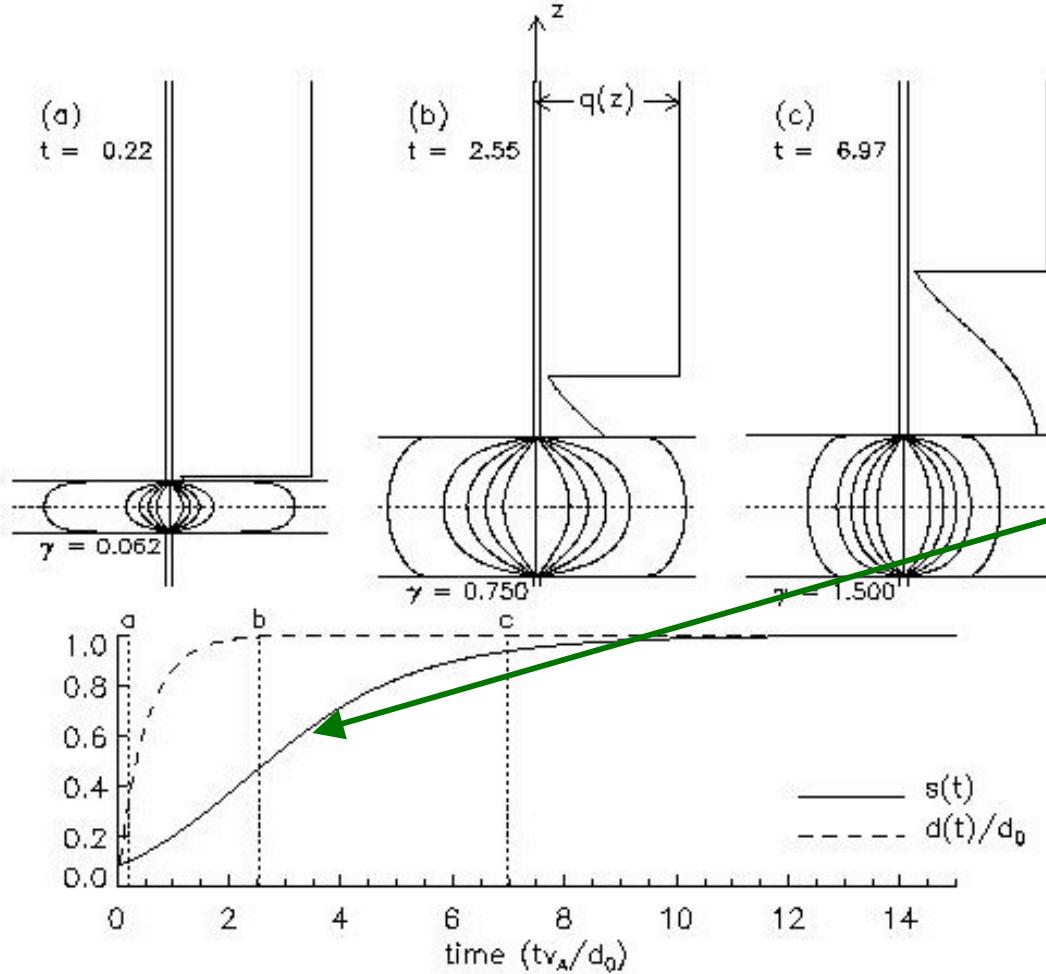
simple model: Longcope & Welsch 1998



- Twist propagates into corona
- Twist-rarefaction waves propagates inward to CZ

Spin from Emergence

simple model: Longcope & Welsch 1998



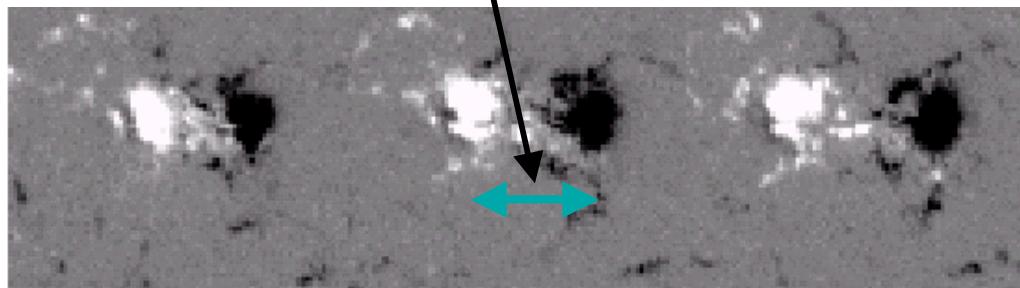
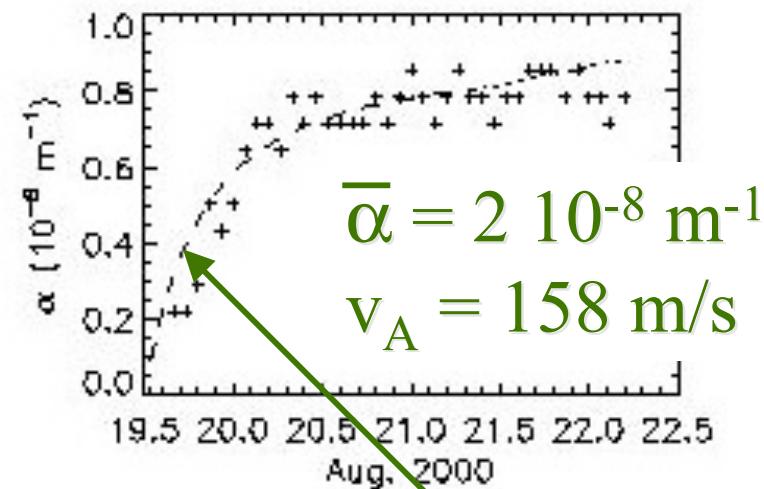
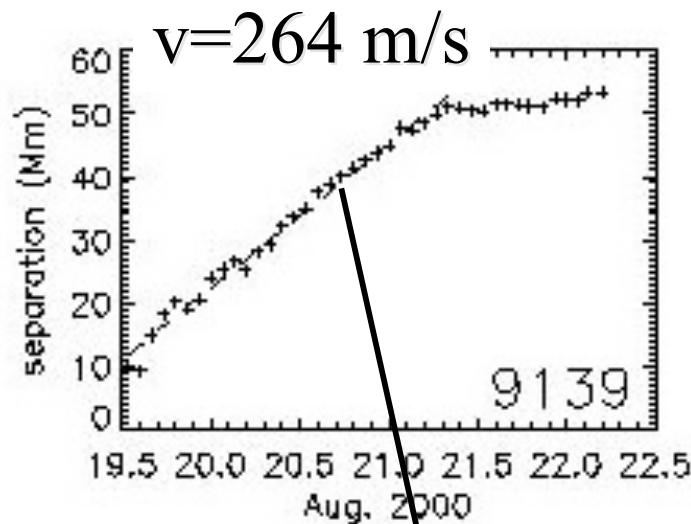
- Twist propagates into corona
- Twist-rarefaction waves propagates inward to CZ
- Characteristic time-scale for adjustment:

$$d/v_A \sim 1 \text{ day}$$

Spin from Emergence

Observation: Pevtsov, Maleev & Longcope 2003

Fit Model to Data

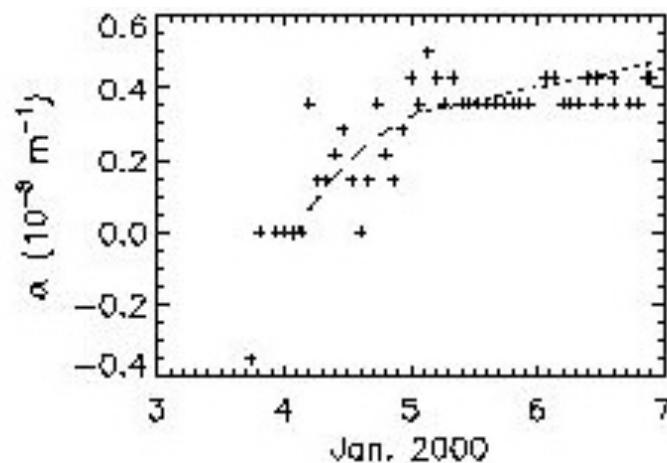
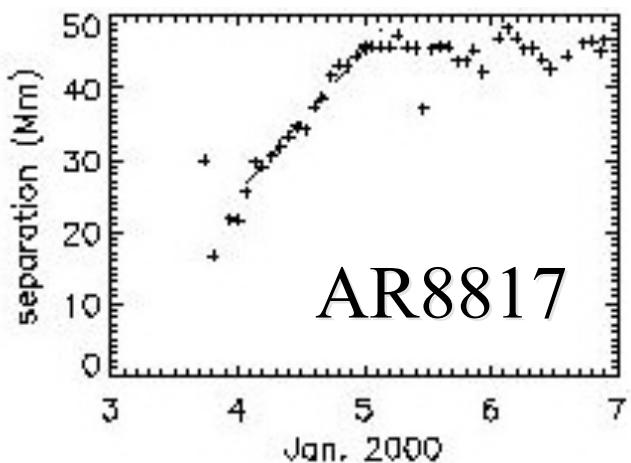
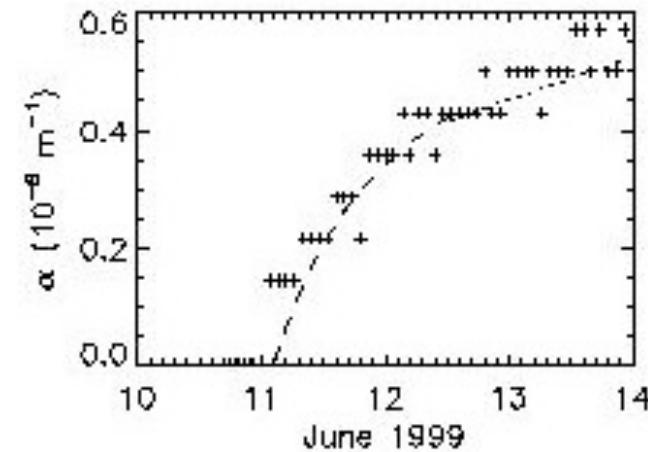
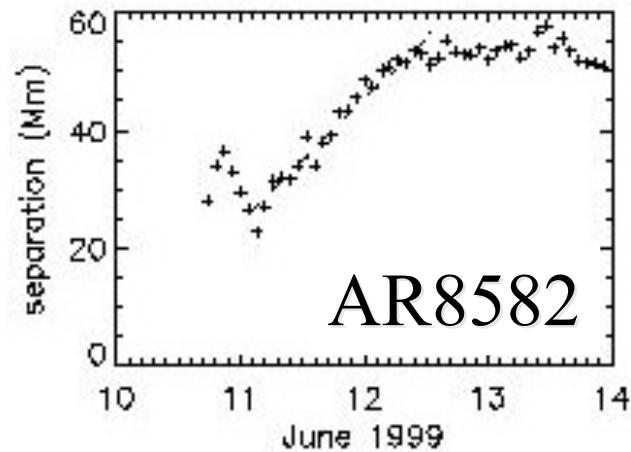


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$$\alpha(t) = \bar{\alpha} \frac{v}{v+1} \left[1 - \left(\frac{d}{d_0} \right)^{-(v+1)} \right]$$

Spin from Emergence

Observation: Pevtsov, Maleev & Longcope 2003

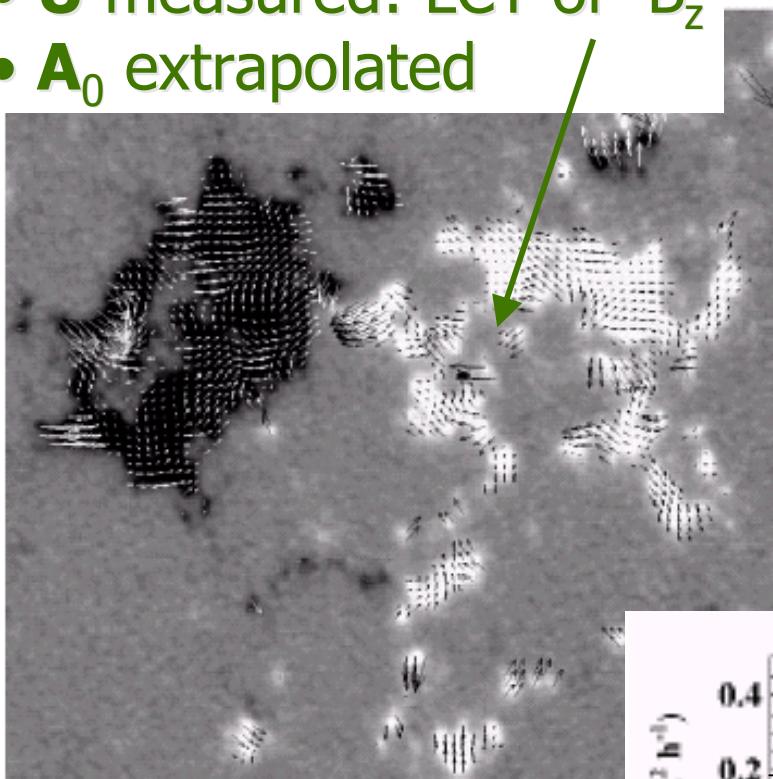


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Measured Velocity

B_z measured: LOS mg

- \mathbf{U} measured: LCT of B_z
- \mathbf{A}_0 extrapolated



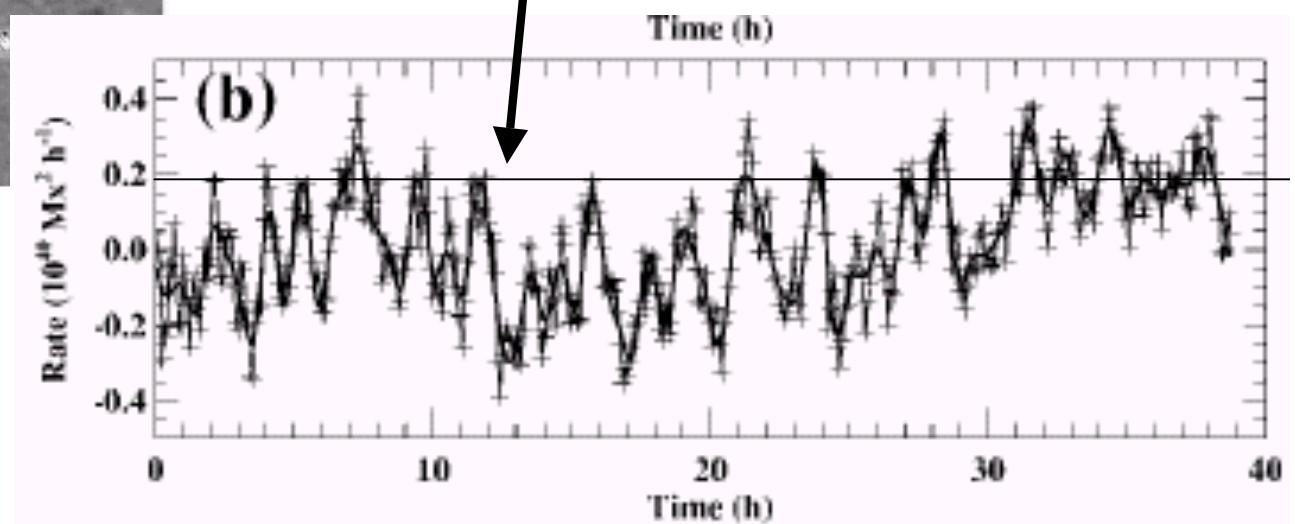
(Chae 2001)

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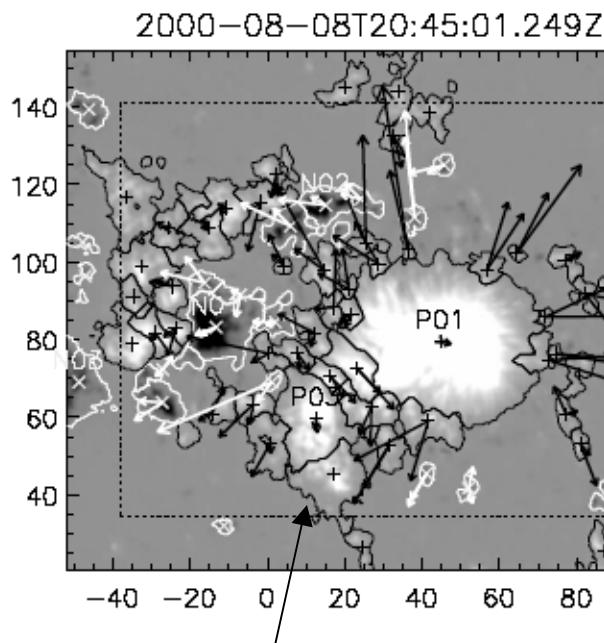
$$\frac{dH_R}{dt} = -2 \int_{z=0} (\mathbf{U} \cdot \mathbf{A}_0) B_z ds$$

$$= +3 \times 10^{40} \text{ Mx}^2/\text{day}$$

$$= \Phi^2 10^{-2}/\text{day}$$

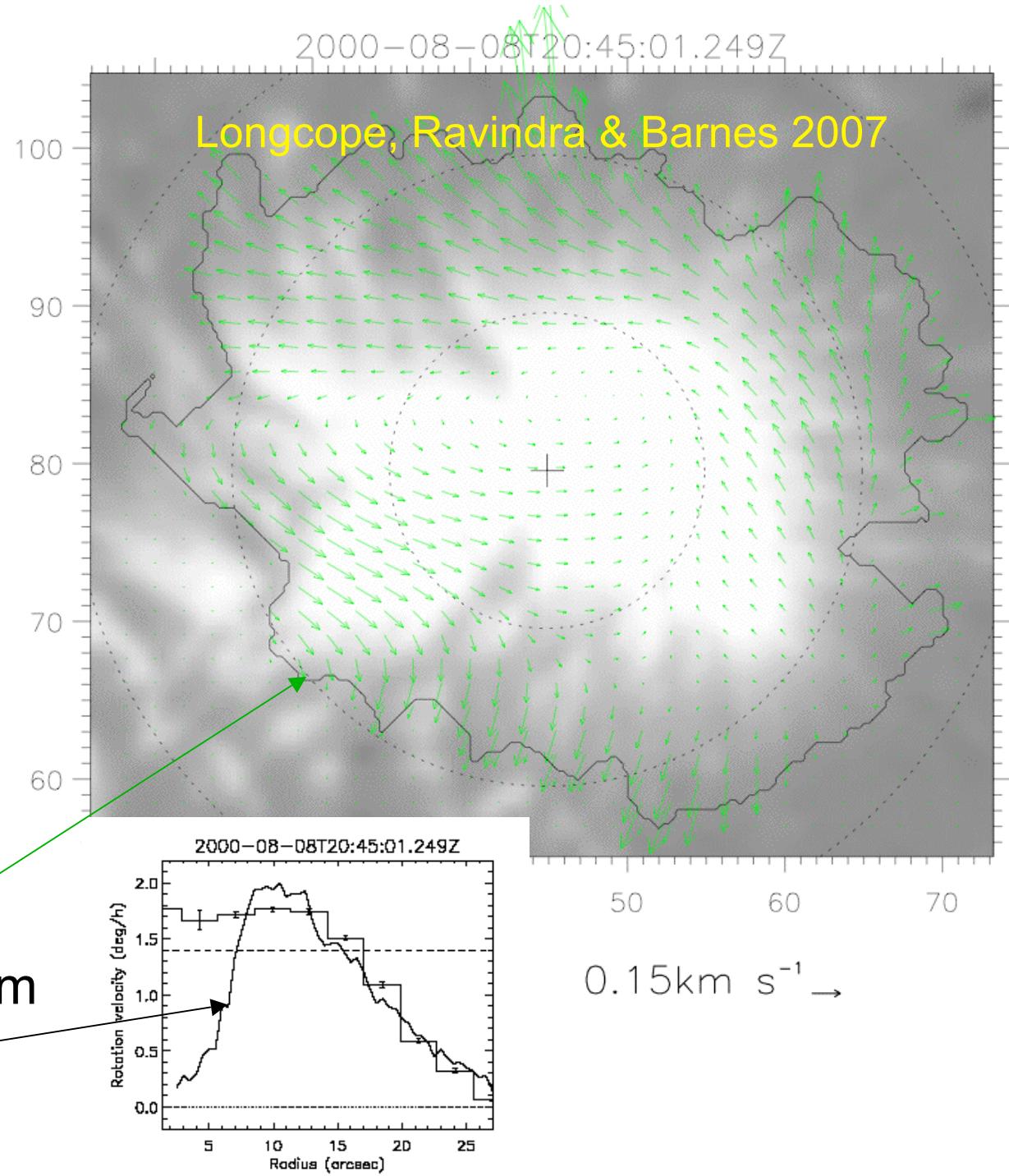


Measuring Spin



- partition m-gram
- $v(x)$ from LCT
- cf WL rotation from Brown et al. 2003

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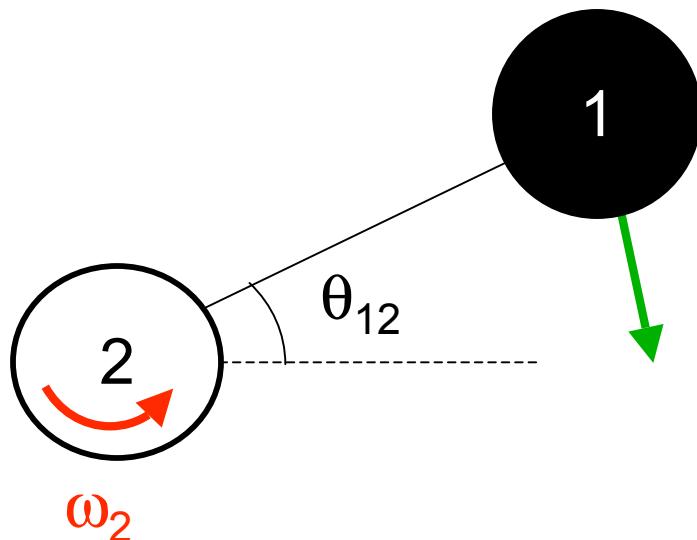
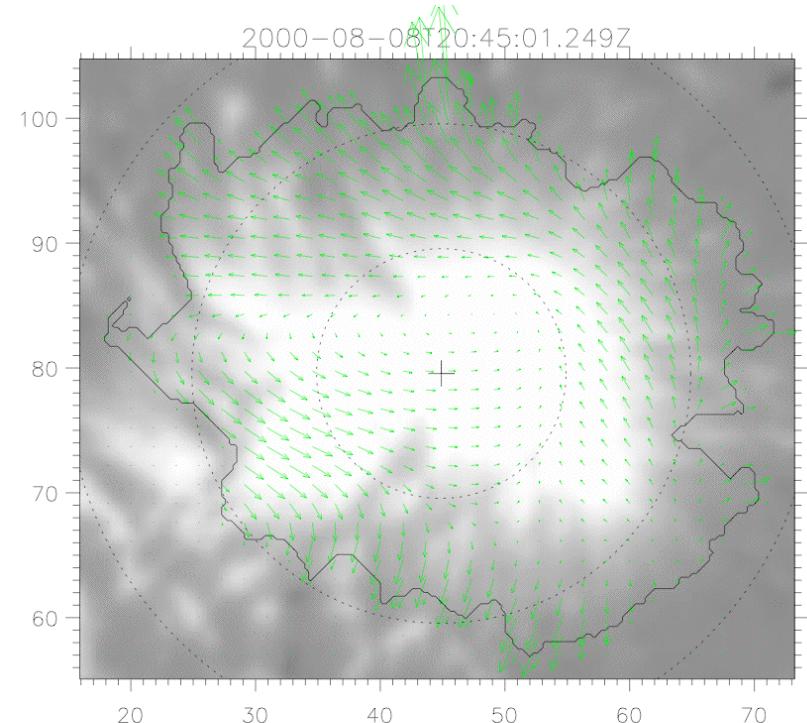


Measuring Spin

$$\frac{dH}{dt} = \sum_a \left(-\frac{\Phi_a^2}{2\pi} \omega_a \right) - \sum_{a,b} \frac{\Phi_a \Phi_b}{2\pi} \frac{d\theta_{ab}}{dt}$$

spin braiding

$$\frac{dH}{dt} = -2 \int_{z=0} (\mathbf{v} \cdot \mathbf{A}_0) B_z dx dy$$



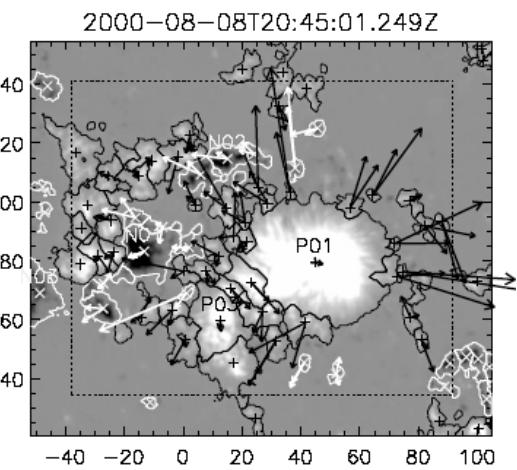
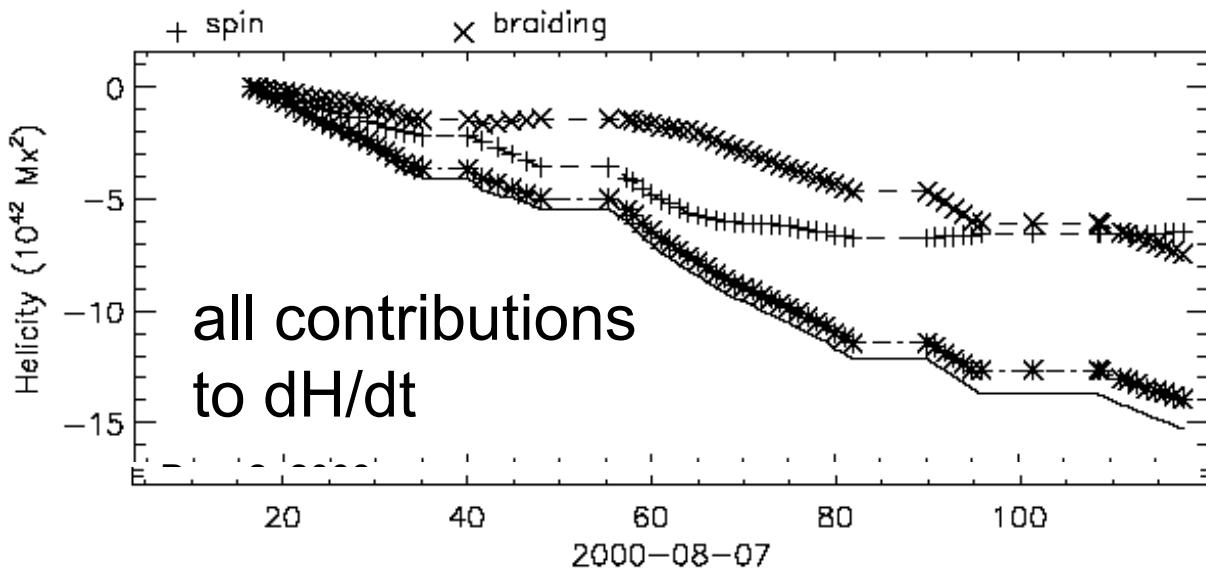
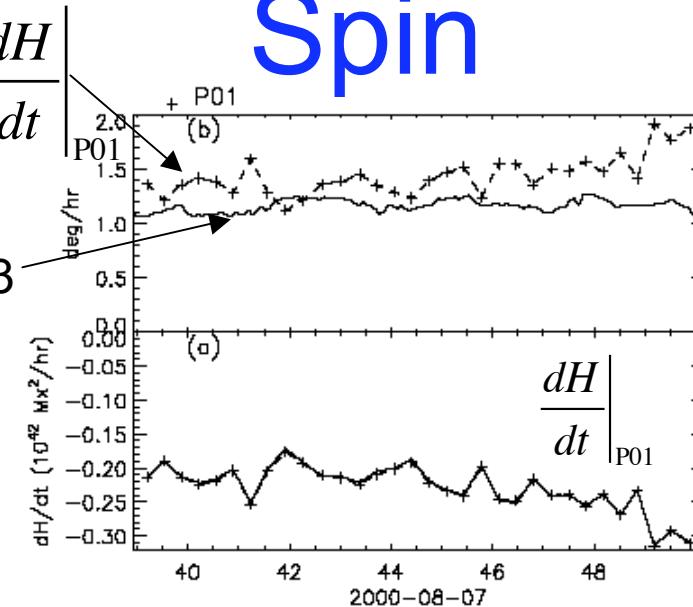
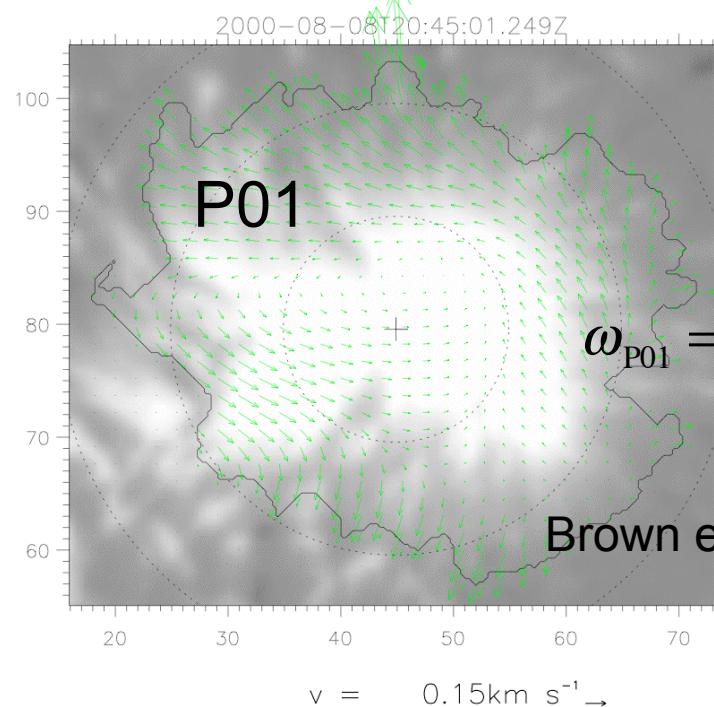
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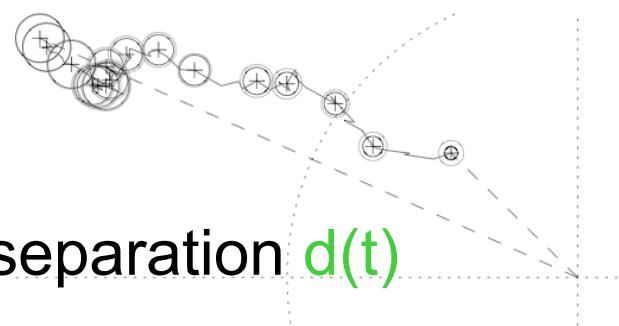
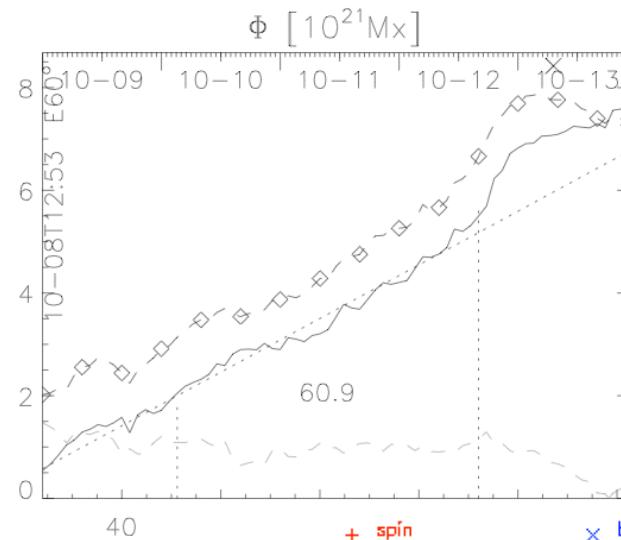
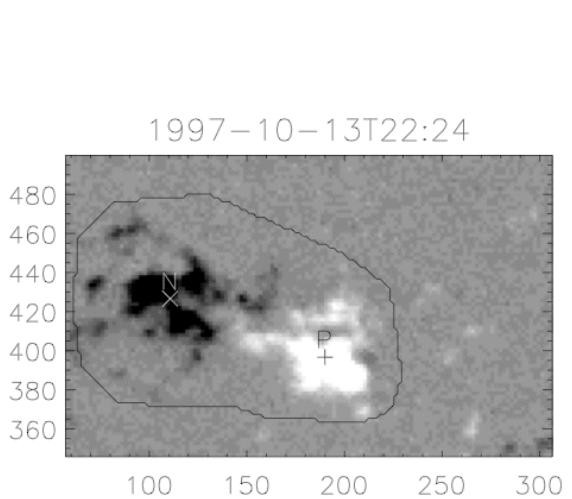
$$\omega_a = - \frac{2\pi}{\Phi_a^2} \frac{dH}{dt} \Big|_a$$

$v = 0.15 \text{ km s}^{-1} \rightarrow$



Measuring Spin



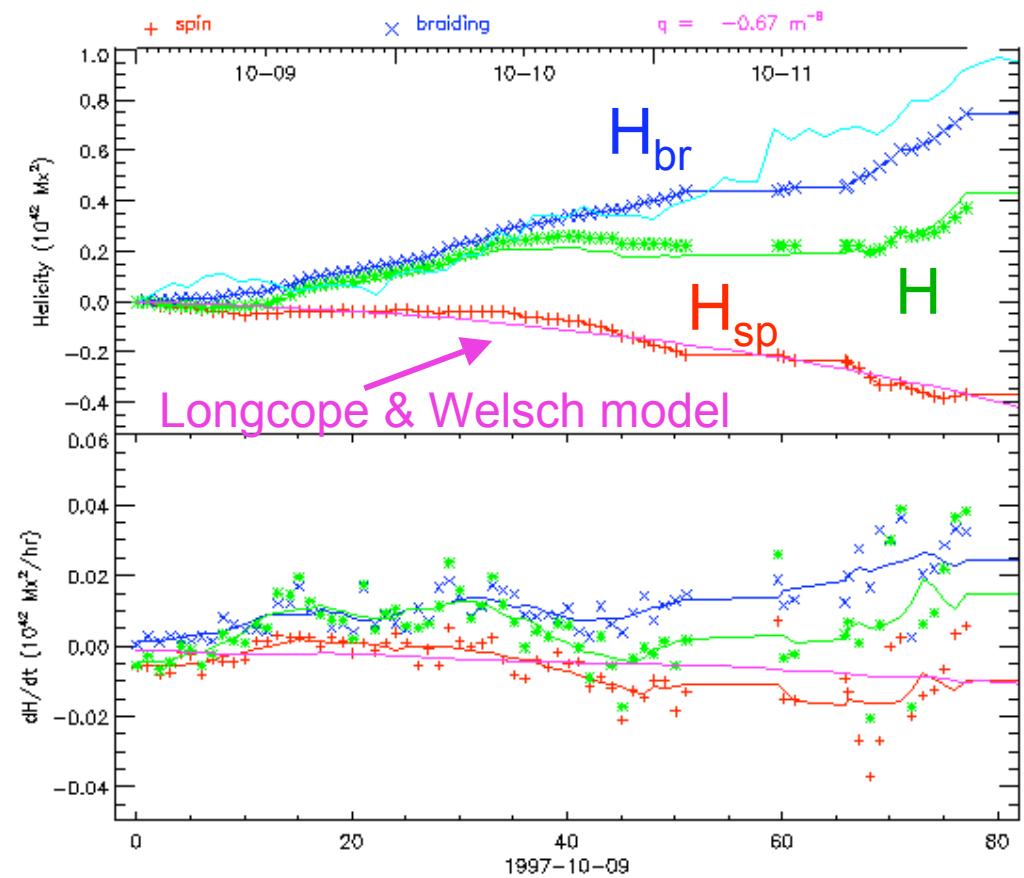


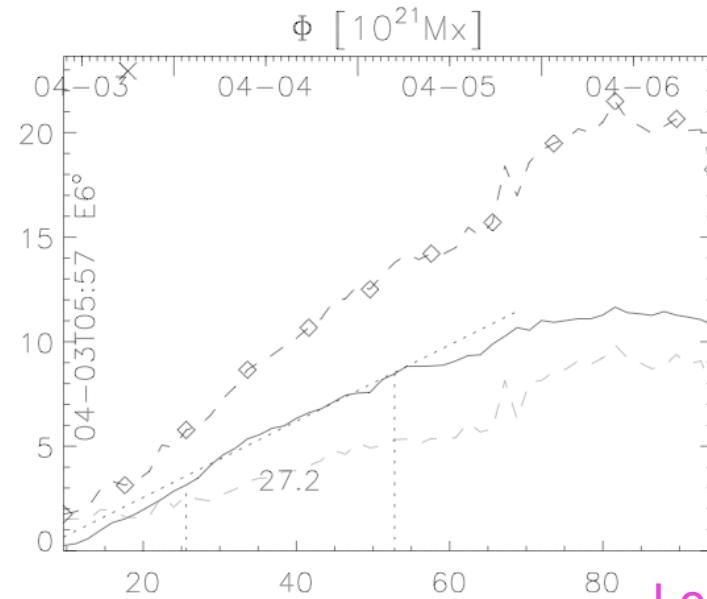
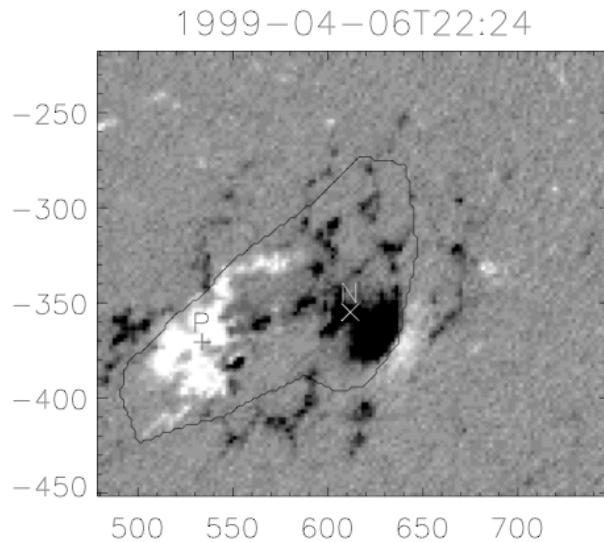
- Helicity dominated by braiding
- Northern AR:
H > 0 H_{sp} < 0

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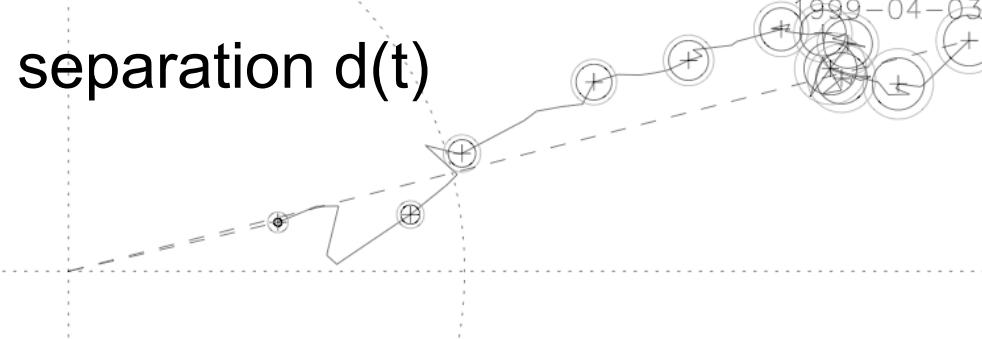
fit:

$$q = -0.67 \times 10^{-8} \text{ m}^{-1}$$

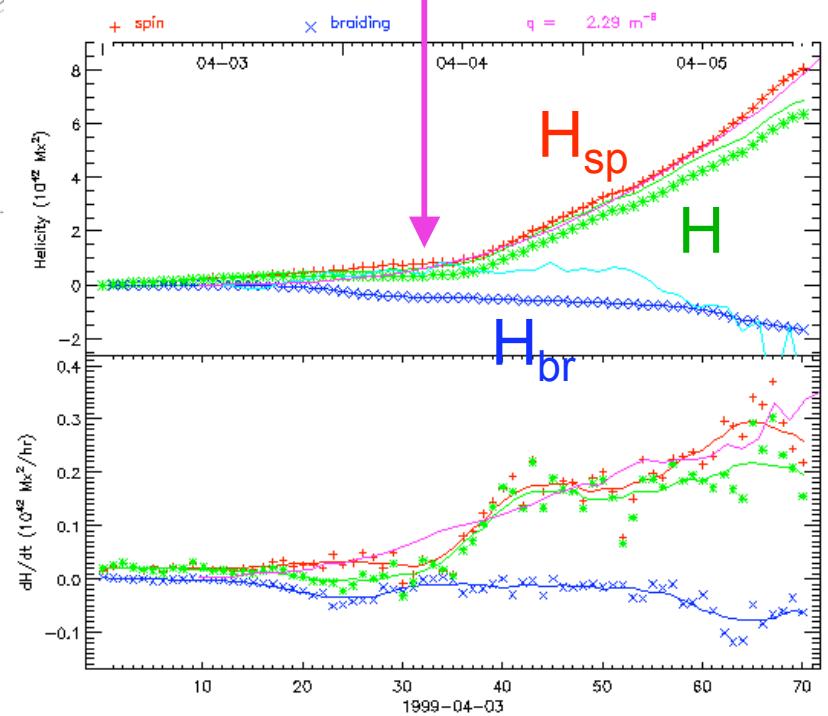




fit:
 $q = +2.3 \times 10^{-8} \text{ m}^{-1}$



LongCOPE & WELSCH model



- Helicity dominated by spin
- Southern AR:

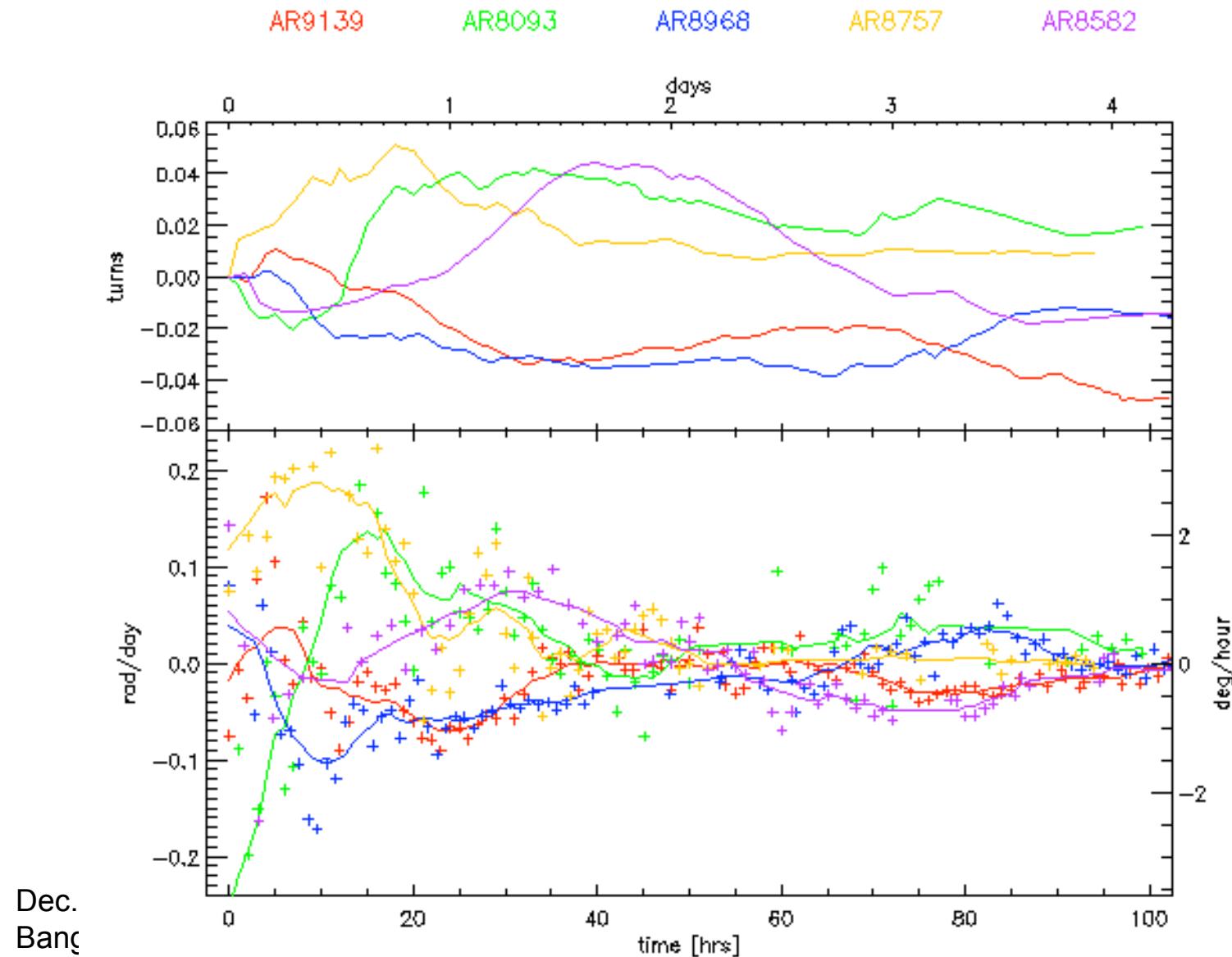
$$H > 0$$

$$H_{sp} > 0$$

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Helicity Injection





Long term helicity injection

AR	Δt (hrs)	$\Delta H/\Phi^2$	$\langle dH/dt \rangle$ (10^{-2} day $^{-1}$)	ref
8011	40	0.003	0.18	Chae (2001)
8100	120	0.02	0.4	Kusano et al. (2002)
8668	50	0.03	1.4	Chae et al. (2001)
9165	80	0.2	6.0	Nindos & Zhang (2002)
10365	120	0.05	1.0	Chae, Moon & Park (2004)
10696	132	0.02	0.4	Lim et al. (2007)

Dec. 2, 2008 (from van Driel-Gesztelyi, Demoulin & Mandrini, 2003)
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Helicity Flux in ARs

$$\frac{dH}{dt} = \Omega \Phi^2 = \frac{\omega}{2\pi} \Phi^2$$

- Differential rotation:
 - Th. (DeVore 2000): $\Omega \sim 3 \times 10^{-3}/\text{day}$
 - Obs. (Demoulin *et al* 2002) $\Omega \sim 3 \times 10^{-4}$
- Proper motions: (observations)
 - LCT (van Driel-Gesztelyi *et al.* 2003) $\Omega \sim 10^{-2}$
 - Sunspot rotation (Brown *et al* 2003) $\Omega \sim 10^{-1}$



Summary

- ARs created by emergence of flux tubes
- Tubes consist of twisted flux -- twisted by turbulence during rising (Σ -effect)
- Helicity of twist propagates into corona
- Observed proper motions (rotating sunspots) consistent with twist propagation