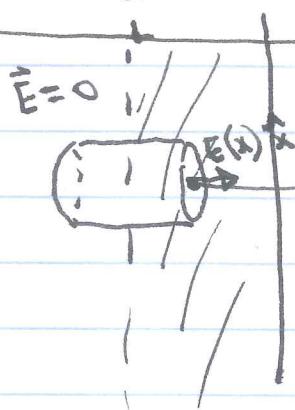


# FORCE ON A CONDUCTOR

§ 2.5.3

1



$$E = \frac{\sigma}{\epsilon_0} \quad (\text{GAUSS'S LAW})$$

$$\sigma(x) = \int_a^x \rho(x') dx', \quad \sigma(b) = \sigma_{\text{tot}}$$

FORCE PER UNIT AREA:

$$\vec{F} = \hat{x} \int_a^b \rho(x) E(x) dx$$

$$= \hat{x} \int_a^b \rho(x) \left( \int_a^x \rho(x') dx' \right) dx$$

Surface =  $\sigma - \int \rho dx$ , integration by parts:  $\rho(x) dx = d\rho$

$$v = \int \rho(x) dx \rightarrow \int_a^x \rho(x) dx'$$

$$u = \int_a^x \rho(x') dx'$$

$$du = \rho(x) dx$$

$$\vec{F} = \frac{\hat{x}}{\epsilon_0} \left[ \left[ \left( \int_a^x \rho(x') dx' \right) \left( \int_a^b \rho(x) dx \right) \Big|_a^b - \int_a^b \rho(x) \left( \int_a^x \rho(x') dx' \right) dx \right] \right]$$

$$\begin{aligned} \frac{\hat{x}}{\epsilon_0} \int_a^b \rho(x) \left( \int_a^x \rho(x') dx' \right) dx &= \frac{\hat{x}}{\epsilon_0} \left[ \left( \int_a^b \rho(x) dx \right) \left( \int_a^b \rho(x') dx' \right) - \int_a^b \rho(x) \left( \int_a^b \rho(x') dx' \right) dx \right] \\ &\quad - \int_a^b \rho(x) \left( \int_a^x \rho(x') dx' \right) dx \end{aligned}$$

$$\therefore 2\vec{F} = \frac{\hat{x}}{\epsilon_0} \left( \int_a^b \rho(x) dx \right)^2 = \frac{\sigma^2 \hat{x}}{\epsilon_0}$$

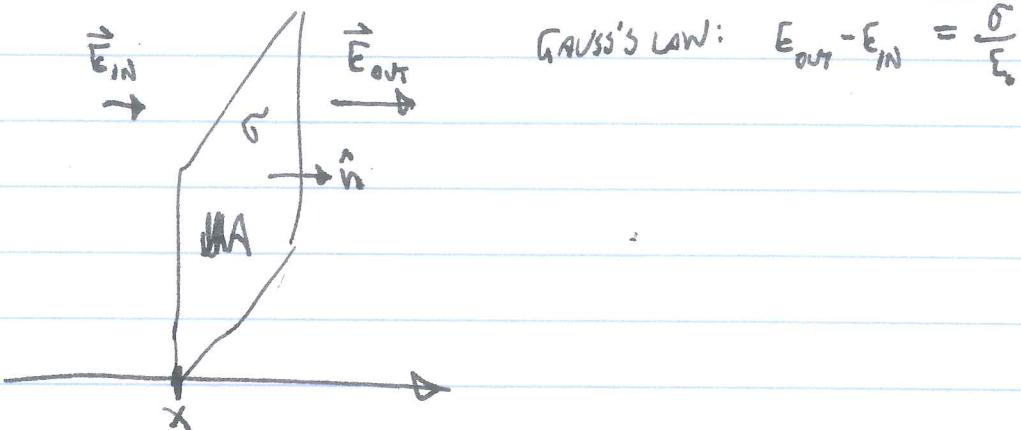
$$\boxed{\vec{F} = \frac{\sigma^2 \hat{x}}{2 \epsilon_0}}$$

THE CHARGE  $\sigma$  INDEPENDENT OF HOW  $x$  IS DISTRIBUTED ALONG THE SURFACE NORMAL.

(2)

## ENERGY ARGUMENT

FORGET CONDUCTORS FOR A MOMENT. THINK OF A SURFACE CHARGE  $\sigma$ ; AREA  $AA$



$$\text{GAUSS'S LAW: } E_{\text{out}} - E_{\text{in}} = \frac{\sigma}{\epsilon_0}$$

~~WE HAVE TO DO THIS~~

HOW MUCH DOES THE ENERGY CHANGE IF I MOVE MY SURFACE  $dx$  TO THE RIGHT?

$$dF \quad dW = \underbrace{AA dx}_{\text{volume}} \left( \frac{1}{2} \epsilon_0 E_{\text{in}}^2 + -\frac{1}{2} \epsilon_0 E_{\text{out}}^2 \right)$$

~~$$\frac{dW}{AA dx} = \frac{1}{2} \epsilon_0 (E_{\text{in}}^2 - E_{\text{out}}^2)$$~~

Q: WHAT IS  $\frac{dW}{AA dx}$ ?

A: A FORCE PER AREA

Q: WHAT DIRECTION IS POSITIVE  $\frac{dW}{AA dx}$ ?

A:  $n-hat$  (To THE <sup>LEFT</sup> ~~RIGHT~~)

LET  $\vec{f} = -\frac{1}{2} \epsilon_0 (E_{\text{in}}^2 - E_{\text{out}}^2) \hat{n} = \frac{1}{2} \epsilon_0 (E_{\text{out}}^2 - E_{\text{in}}^2) \hat{n}$

THIS FORCE IS AN ELECTROSTATIC TENSION. THE CHARGE IS "PULLED" INTO THE STRONGER FIELD REGION.

CONDUCTOR:  $E_{\text{in}} = 0$ , so  $\boxed{\vec{f} = \frac{1}{2} \epsilon_0 E_{\text{out}}^2 \hat{n}}$

Q: ~~WHAT'S~~ CAN A CONDUCTOR "FEEL" A FORCE IN  $-\hat{n}$  DIRECTION?  
A: NO!

APPLYING OUR RESULT TO A CONDUCTOR,  
DO OUR RESULTS MATCH?

$$E_{\text{out}} = \frac{\sigma}{\epsilon_0}, \text{ so } \underline{\text{YES}},$$

$$\vec{F} = \frac{1}{2} \epsilon_0 E^2 \hat{n} = \cancel{\mu_0} \frac{\sigma^2 \hat{n}}{2 \epsilon_0}$$

Force on a conductor

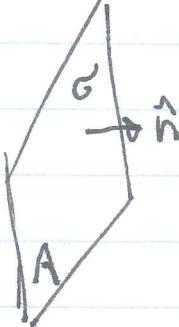
### THIRD DERIVATION

NOTICE THAT WE COULD ALSO WRITE

$$\vec{F} = \frac{1}{2} \sigma E \hat{n}.$$

IT IS AS IF ONLY HALF THE CHARGE EXPERIENCES  $E$ , OR  
ONLY HALF OF  $\epsilon$  INFLUENCES THE CHARGE. THIS IS  
ACTUALLY EASY TO UNDERSTAND.

IMAGINE AN ISOLATED SURFACE CHARGE  $\sigma$ :

$$\vec{E}_{\text{in}}^I = \frac{\sigma \hat{n}}{2 \epsilon_0}$$


$$\vec{E}_{\text{out}}^I = \frac{\sigma \hat{n}}{2 \epsilon_0}$$

SINCE  $\sigma$  IS ISOLATED,  
AND IT CAN'T EXERT  
A FORCE ON ITSELF,  
 $\vec{F}^I = 0$ .

NOW IMPOSE AN EXTERNAL FIELD,  $\vec{E}_{\text{ext}} = E_{\text{ext}} \hat{n}$ .

THE FORCE IS CLEARLY  $\vec{F} = \sigma \vec{E}_{\text{ext}}$

NOW IN THE CASE OF THE CONDUCTOR,

$$\vec{E}_{\text{in}} = \vec{E}_{\text{in}}^I + \vec{E}_{\text{ext}} = 0 \Rightarrow \vec{E}_{\text{ext}} = \frac{\sigma \hat{n}}{2 \epsilon_0}$$

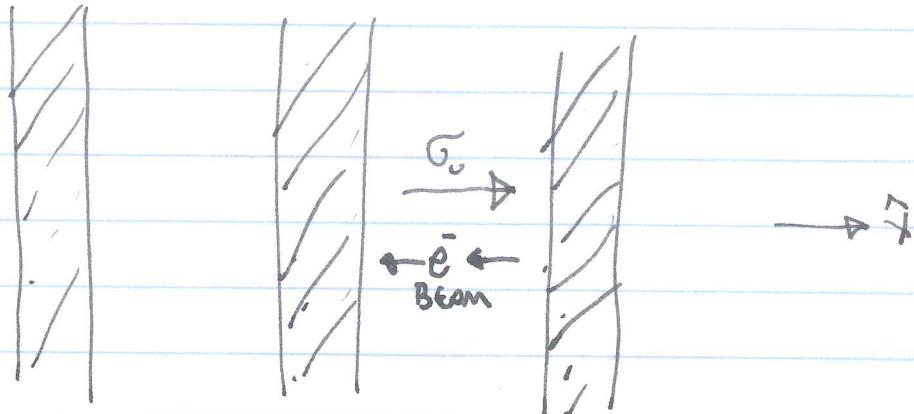
$$\therefore \boxed{\vec{F} = \frac{\sigma^2 \hat{n}}{2 \epsilon_0}} \quad \checkmark$$

THIS DERIVATION IS MORE SIMILAR TO THAT IN GRIFFITHS.

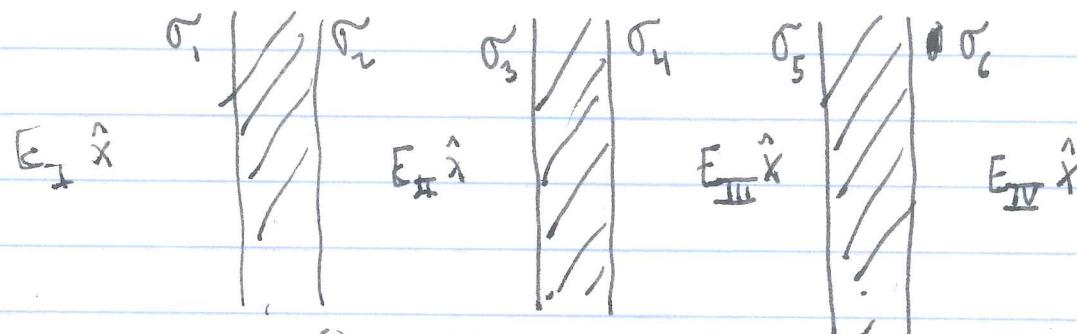
4

## EXAMPLE

3 CONDUCTING PLATES, ALL NEUTRAL,  $V = \sigma = 0$  EVERYWHERE.



NOW A CHARGE  $G_0$  IS TRANSFERRED FROM THE MIDDLE PLATE TO THE RIGHT PLATE. HOW DOES THE CHARGE DISTRIBUTE ITSELF? WHAT IS THE FORCE ON EACH PLATE? WHAT ARE THE FIELDS?



$$\text{BUT } \sigma_1 + \sigma_2 = 0 \quad (1) \quad \sigma_3 + \sigma_4 = -\sigma_0 \quad (2) \quad \sigma_5 + \sigma_6 = \sigma_0 \quad (3)$$

Q: WHAT ELSE DO I KNOW?

A:  $\vec{E} = 0$  IN ALL 3 CONDUCTORS!  $\Rightarrow$  3 MORE EQUATIONS:

$$(4) \quad \sigma_1 - (\underbrace{\sigma_2 + \sigma_3 + \sigma_4 + \sigma_5 + \sigma_6}_{=0}) = 0 \Rightarrow \sigma_1 - 0 = 0 \Rightarrow \boxed{\sigma_1 = \sigma_2 = 0}$$

$$(5) \quad \cancel{\sigma_1 + \sigma_2} + \sigma_3 - (\sigma_4 + \sigma_5 + \sigma_6) = 0 \Rightarrow -\sigma_0 - 2\sigma_4 + \sigma_0 = 0 \therefore \boxed{\sigma_4 = 0}$$

$$\boxed{\sigma_3 = 0}$$

$$(6) \quad \sigma_1 + \sigma_2 + \sigma_3 + \sigma_4 + \sigma_5 - \sigma_6 = 0 \Rightarrow -\sigma_0 + (\sigma_5 + \sigma_6) \cdot 2\sigma_4 = 0 \\ \Rightarrow \boxed{\sigma_6 = 0} \quad \boxed{\sigma_5 = \sigma_0}$$

(5)

WHAT ARE THE ~~FORCES~~?  
FIELDS?

$$E_I = 0, \quad E_{II} = 0, \quad E_{III} = \frac{\sigma_4 \hat{x}}{2\epsilon_0} = -\frac{\sigma_5 \hat{x}}{2\epsilon_0} = -\frac{\sigma_6 \hat{x}}{2\epsilon_0}$$

FORCES:  $\vec{F}_L = 0$

LEFT PLATE:  $\sigma_1 = \sigma_2 = 0 \therefore \vec{F}_L = 0$

MIDDLE PLATE:

$$\vec{F}_M = \frac{\sigma_4 \hat{x}}{2\epsilon_0} = \frac{\sigma_6 \hat{x}}{2\epsilon_0}$$

~~RIGHT~~  
~~LEFT PLATE~~

RIGHT PLATE:

$$\vec{F}_R = -\frac{\sigma_5 \hat{x}}{2\epsilon_0} = -\frac{\sigma_6 \hat{x}}{2\epsilon_0}$$

NOTE THAT  $\vec{F}_{NET} = \vec{F}_L + \vec{F}_M + \vec{F}_R = 0$ .

Q: DID THAT HAVE TO BE TRUE?

A: YES! NET FORCE ON AN ISOLATED SYSTEM IS ALWAYS ZERO.

Q: WHAT IF THE "ELECTRON BEAM" WENT THROUGH A TINY HOLE IN THE MIDDLE PLATE AND LANDED ON THE LEFT PLATE INSTEAD?