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SO FAR IN CH. 21, PDES

- SEPARATION OF VARIABLES
- INTEGRAL TRANSFORMS
- GREEN'S FUNCTIONS FOR HOMOGENEOUS B.C.'S
 - METHOD OF IMAGES
 - METHOD OF EIGENFUNCTION EXPANSION

TODAY — § 21.5.2, GREEN'S FUNCTIONS FOR
INHOMOGENEOUS B.C.'S

IT'S A THEORY DAY — NO EXAMPLES!

UPCOMING — § 21.5.3 G.F. FOR DIRICHLET B.C.'S

§ 21.5.4 G.F. FOR NEUMANN B.C.'S

- AND THAT'S THE END OF THE CHAPTER!
- KNOWING HOW TO SOLVE INHOMOGENEOUS PDES WITH INHOMOGENEOUS B.C.'S IS THE PINNACLE OF THE THEORY FOR LINEAR PDES.
- YOU MAY FIND IN THE RECENT LITERATURE G.F. SOLUTIONS OF THIS TYPE FOR PROBLEMS SUCH AS

$$\nabla \times \vec{B} = \alpha \vec{B},$$

WITH $\vec{B} \cdot \hat{n}$ GIVEN ON THE BOUNDARY.

PDE GREEN'S FUNCTIONS FOR INHOMOGENEOUS B.C.'S

(§ 21.5.2)

WE'LL FOCUS ON POISSON'S EQUATION, AS DOES RHB.

TO BE SOLVED:

- $\nabla^2 u(\vec{r}) = \rho(\vec{r})$, $\vec{r} \in V$ BOUNDED BY S'
- INHOMOGENEOUS NEUMANN OR DIRICHLET B.C.'S ON S'
- THE GENERAL APPROACH WORKS FOR OTHER OPERATORS, BUT WE FOCUS ON ∇^2 .
- MY PRESENTATION DIFFERS SOMEWHAT FROM RHB.

BEGIN BY SOLVING

$$\nabla^2 G(\vec{r}, \vec{r}_0) = \delta(\vec{r} - \vec{r}_0); \quad \vec{r}, \vec{r}_0 \in V \quad \star\star$$

- DO NOT SUBJECT G TO THE B.C.'S
 → G IS THEREFORE NOT UNIQUE.

START WITH THE SUPERPOSITION

$$\int_V G(\vec{r}, \vec{r}_0) \rho(\vec{r}_0) dV_0 \quad \neq \underline{\bar{u}(\vec{r})}$$

- DOESN'T SATISFY B.C.'S, EVEN IF G DOES,
 BECAUSE B.C.'S ARE INHOMOGENEOUS.
- DOES SATISFY THE PDE $\nabla^2 \bar{u} = \rho$. (\star)

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NOW SEEK A CORRECTION TERM $h(\vec{r})$ SUCH THAT

$$u(\vec{r}) = \int_V G(\vec{r}, \vec{r}_0) \rho(\vec{r}_0) dV_0 + h(\vec{r}) \quad (*)$$

Q: SPECULATE: HOW CAN WE CHOOSE G AND h ?

A: ① LET G SATISFY HOMOGENEOUS B.C.'S, SO $\int G \rho dV_0$ DOES TOO.

② LET $h(\vec{r})$ SATISFY $\nabla^2 h = 0$

③ LET $h(\vec{r})$ SATISFY THE INHOMOGENEOUS B.C.'S FOR $u(\vec{r})$.

WE ASSUME NONE OF THIS, BUT IT WILL FALL OUT...

FIRST, EXPRESS $u(\vec{r})$ (TRIVIALLY!) AS

$$u(\vec{r}) = \int_V u(\vec{r}_0) \delta(\vec{r} - \vec{r}_0) dV_0$$

SO THAT IT CAN BE PULLED INTO THE INTEGRAL
WITH G . \otimes BECOMES:

$$\begin{aligned} h(\vec{r}) &= \int_V \left[-G(\vec{r}, \vec{r}_0) \rho(\vec{r}_0) + u(\vec{r}_0) \delta(\vec{r} - \vec{r}_0) \right] dV_0 \\ &= \int_V \left[-G(\vec{r}, \vec{r}_0) \overset{\downarrow \star}{\nabla^2} u(\vec{r}_0) + u(\vec{r}_0) \overset{\downarrow \star \star}{\nabla^2} G(\vec{r}, \vec{r}_0) \right] dV_0 \end{aligned}$$

STRATEGY

IF WE CAN CHANGE THIS ∇^2 TO ∇_0^2 , THEN
WE'LL HAVE THE LHS OF GREEN'S 2nd THEOREM.
THIS WILL GIVE US A SURFACE INTEGRAL (ON S'),
WHICH WILL CONNECT US TO THE B.C.'S.



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TAKE \star AND SWITCH VARIABLES, $\vec{r} \leftrightarrow \vec{r}_0$:

$$\nabla_{\vec{r}_0}^2 G(\vec{r}_0, \vec{r}) = \delta(\vec{r}_0 - \vec{r})$$

BUT δ IS SYMMETRIC ABOUT THE ORIGIN, SO

$$\nabla_{\vec{r}_0}^2 G(\vec{r}_0, \vec{r}) = \delta(\vec{r} - \vec{r}_0) = \nabla_{\vec{r}}^2 G(\vec{r}, \vec{r}_0)$$

SUPPOSE G SATISFIES HOMOGENEOUS B.C.'S, SO THAT
 ∇^2 IS HERMITIAN IN \star (THOUGH NOT IN \star !)

$$\Rightarrow G(\vec{r}_0, \vec{r}) = G^*(\vec{r}, \vec{r}_0) \quad (\text{RESULT OF LAST LECTURE})$$

AND THUS

$$\nabla_{\vec{r}_0}^2 G^*(\vec{r}, \vec{r}_0) = \nabla_{\vec{r}}^2 G(\vec{r}, \vec{r}_0)$$

ASSUME FURTHER $G \in \mathbb{R}$ (TRUE IF THE OPERATOR,
 $\rho(\vec{r})$, AND B.C.'S $\in \mathbb{R}$). THEN WE HAVE

$$\nabla_{\vec{r}_0}^2 G(\vec{r}, \vec{r}_0) = \nabla_{\vec{r}}^2 G(\vec{r}, \vec{r}_0) \quad !!!$$

SO THEN

$$h(\vec{r}) = \int_V \left[u(\vec{r}_0) \nabla_{\vec{r}_0}^2 G(\vec{r}, \vec{r}_0) - G(\vec{r}, \vec{r}_0) \nabla_{\vec{r}_0}^2 u(\vec{r}_0) \right] dV_0$$

N.B. : OUR PRESENTATION SO FAR IS GENERIC,
 UP TO THIS POINT, $\nabla_{\vec{r}_0}^2$ COULD BE ANY
REAL OPERATOR \mathcal{D} , WITH G ANY GREEN'S
 FUNCTION FOR \mathcal{D} , SATISFYING HOMOGENEOUS
 B.C.'S, SO LONG AS IT BECOMES HERMITIAN
 WHEN HOMOGENEOUS B.C.'S ARE APPLIED.

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Now, we get specific to ∇^2 . Using GREEN'S SECOND THEOREM,

$$h(\vec{r}) = \oint_S [u(\vec{r}_0) \nabla_0 \cdot G(\vec{r}, \vec{r}_0) - G(\vec{r}, \vec{r}_0) \nabla_0 u(\vec{r}_0)] \cdot \hat{n}_0 dA_0$$

$$= \oint_S [u(\vec{r}_0) \frac{\partial G(\vec{r}, \vec{r}_0)}{\partial n_0} - G(\vec{r}, \vec{r}_0) \frac{\partial u(\vec{r}_0)}{\partial n_0}] dA_0$$

Therefore, going back to \otimes ,

$$u(\vec{r}) = \int_V G(\vec{r}, \vec{r}_0) \rho(\vec{r}_0) dV_0 + \oint_S [u(\vec{r}_0) \frac{\partial G(\vec{r}, \vec{r}_0)}{\partial n_0} - G(\vec{r}, \vec{r}_0) \frac{\partial u(\vec{r}_0)}{\partial n_0}] dA_0$$

Cf. Eq. (21.84).

Before we go further, recall that we had made 3 predictions about G & h to solve this problem:

- ① Let G satisfy homogeneous B.C.'s \leftarrow This was used in our derivation of $h(\vec{r})$, but there are other approaches.
- ② Let $h(\vec{r})$ satisfy $\nabla^2 h = 0$ on V (inside S .)
- ③ Let $h(\vec{r})$ satisfy the inhomogeneous B.C.'s.

Q: Can we agree that these conditions logically entail \otimes ? Discuss.

A: Yes, they do.



(5)

CASE I : DIRICHLET B.C.'S, $u(\vec{r}) = \phi(\vec{r})$ FOR $\vec{r} \in S$

① LET $G(\vec{r}, \vec{r}_0) = 0$ FOR $\vec{r}_0 \in S$ (WHY? LOOK AT 21.84).

IT ALSO FOLLOWS THAT $G = 0$ FOR $\vec{r} \in S$ (WHY?)

$$\text{THEN } h(\vec{r}) = \oint_S [u(\vec{r}_0) \frac{\partial G(\vec{r}, \vec{r}_0)}{\partial n_0}] dA_0$$

MIGHT BE PEDAGOGICALLY

USEFUL TO TAKE IT ANOTHER
WAY: $\frac{\partial}{\partial n_0} g(\vec{r} - \vec{r}_0)$

$$\left. \begin{aligned} ② \nabla^2 h &= \oint_S [u(\vec{r}_0) \frac{\partial}{\partial n_0} \nabla^2 G(\vec{r}, \vec{r}_0)] dA_0 \\ &= \oint_S [u(\vec{r}_0) \frac{\partial}{\partial n_0} \delta(\vec{r} - \vec{r}_0)] dA_0 \end{aligned} \right\}$$

• NOTE THAT $g(\vec{r} - \vec{r}_0) = 0$ FOR $\vec{r} \neq \vec{r}_0$.

• THUS, $\nabla^2 h(\vec{r}) = 0$ FOR \vec{r} INSIDE (NOT ON) S . ✓

③ DOES $h(\vec{r}) = \phi(\vec{r})$ FOR $\vec{r} \in S$?

$$\begin{aligned} h(\vec{r}) &= \oint_S [\phi(\vec{r}_0) \nabla_{\vec{r}_0} G(\vec{r}, \vec{r}_0)] \cdot \hat{n} dA_0 \\ &= \int_V \nabla_{\vec{r}_0} [\phi(\vec{r}_0) \nabla_{\vec{r}_0} G(\vec{r}, \vec{r}_0)] dV_0 \quad (\text{GAUSS'S THM}) \\ &= \int_V \left\{ [\nabla_{\vec{r}_0} \phi(\vec{r}_0)] \cdot [\nabla_{\vec{r}_0} G(\vec{r}, \vec{r}_0)] \right. \\ &\quad \left. + \phi(\vec{r}_0) \underbrace{\nabla_{\vec{r}_0}^2 G(\vec{r}, \vec{r}_0)}_{\delta(\vec{r} - \vec{r}_0)} \right\} dV_0 \\ &= \phi(\vec{r}) + \underbrace{\int_V [\nabla_{\vec{r}_0} \phi(\vec{r}_0)] \cdot [\nabla_{\vec{r}_0} G(\vec{r}, \vec{r}_0)]}_{\uparrow} \end{aligned}$$

Q: WHAT ABOUT THE SECOND TERM? ↗

A: SINCE $\phi(\vec{r}_0)$ IS DEFINED ONLY FOR $\vec{r}_0 \in S$,
WE CAN ONLY EVALUATE THE TANGENTIAL ($\perp \vec{n}_0$)
COMPONENTS, $\nabla_0 \phi(\vec{r}_0) \rightarrow \nabla_0^{\parallel} \phi(\vec{r}_0)$. NOW,

$$\nabla_0^{\parallel} G(\vec{r}, \vec{r}_0) = 0 \text{ BECAUSE OF THE B.C.}$$

(G DOESN'T VARY OVER THE SURFACE!)

SO $[\nabla_0 \phi(\vec{r}_0)] \cdot [\nabla_0 G(\vec{r}, \vec{r}_0)] = 0 \text{ ON } S.$

$$\therefore h(\vec{r}) = \phi(\vec{r}) \text{ ON } S. \quad \checkmark$$

FOR AN ADDED EXERCISE, YOU MAY SHOW THAT PROPOSITIONS
①, ② & ③ ARE SATISFIED WHEN THE B.C.'S ARE NEUMANN.

GENERAL RULE: (§§ 21.5.3 - 21.5.4)

- FOR DIRICHLET CONDITIONS $u(\vec{r}) = \phi(\vec{r})$, $\vec{r} \in S$,
USE $G(\vec{r}, \vec{r}_0) = 0$ FOR $\vec{r}_0 \in S$

- FOR NEUMANN CONDITIONS $\frac{\partial}{\partial n} u(\vec{r}) = \psi(\vec{r})$, $\vec{r} \in S$,
USE $\frac{\partial}{\partial n} G(\vec{r}, \vec{r}_0) = 0$ FOR $\vec{r}_0 \in S$

Q: COULD WE HAVE ANTICIPATED THIS BASED ON